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Factorised 3d $\mathcal{N} = 4$ orthosymplectic quivers

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ABSTRACT: We study the moduli space of 3d $\mathcal{N} = 4$ quiver gauge theories with unitary, orthogonal and symplectic gauge nodes, that fall into exceptional sequences. We find that both the Higgs and Coulomb branches of the moduli space factorise into decoupled sectors. Each decoupled sector is described by a single quiver gauge theory with only unitary gauge nodes. The orthosymplectic quivers serve as magnetic quivers for 5d $\mathcal{N} = 1$ superconformal field theories which can be engineered in type IIB string theories both with and without an O5 plane. We use this point of view to postulate the dual pairs of unitary and orthosymplectic quivers by deriving them as magnetic quivers of the 5d theory. We use this correspondence to conjecture exact highest weight generating functions for the Coulomb branch Hilbert series of the orthosymplectic quivers, and provide tests of these results by directly computing the Hilbert series for the orthosymplectic quivers in a series expansion.

KEYWORDS: Brane Dynamics in Gauge Theories, Supersymmetric Gauge Theory, Field Theories in Higher Dimensions, Field Theories in Lower Dimensions

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1 Introduction and summary of results

Gauge theories in three spacetime dimensions are strongly coupled in the IR, determining their low energy dynamics is therefore generically difficult. One arena in which one can overcome this difficulty is the realm of 3d $\mathcal{N} = 4$ gauge theories. Their relevance to string theory was highlighted very early after the D-brane revolution in a landmark paper by Hanany and Witten [1], which facilitated further explorations of the subject. A more

recent development is to use 3d $\mathcal{N} = 4$ theories as a probe to study higher dimensional superconformal field theories (SCFTs) as well as gauge theories [2–14]. The idea is to relate the Higgs branch of these higher dimensional theories to the Coulomb branch of the 3d theory, the latter of which is dubbed magnetic quiver (MQ). In addition to their significance to string theory or higher dimensional theories, 3d $\mathcal{N} = 4$ theories possess rich dynamics, making them interesting objects in their own right. Notable among their many rich properties is mirror symmetry [1, 15, 16], a duality which relates pairs of 3d $\mathcal{N} = 4$ gauge theories where the role of masses and Fayet-Iliopoulos terms are exchanged. The moduli space of vacua of a 3d $\mathcal{N} = 4$ gauge theory is generically comprised of a Higgs branch and a Coulomb branch,¹ which are exchanged under mirror symmetry. In certain cases, it can be shown that the magnetic quiver of a given theory, is the 3d mirror of the torus compactification of that theory to 3d [7, 17–19]. Another interesting aspect of 3d $\mathcal{N} = 4$ theories is the enhancement of their global symmetry in the infrared (IR) limit. Gauge theories in 3d possess a topological (or magnetic) symmetry, which is valued in the centre of the Langland dual of the gauge group G . Classically, this is an abelian symmetry, but in the IR this is typically enhanced to a non-abelian global symmetry.

It will be convenient for us to make a distinction between quivers which are made entirely of unitary gauge nodes, and those which can also have orthogonal and/or symplectic gauge nodes in addition. We will refer to the former as unitary and the latter as orthosymplectic (OSp) quivers respectively. The primary focus of this paper is orthosymplectic magnetic quivers.

A convenient tool to study the moduli space of 3d $\mathcal{N} = 4$ theories is the Hilbert series, which enumerates gauge invariant operators graded by their conformal dimension. The Coulomb branch Hilbert series can be computed using the monopole formula [20], while the Higgs branch Hilbert series can be evaluated using the Molien-Weyl formula [21]. The Coulomb branch Hilbert series is sensitive to the pattern of symmetry enhancement discussed above, so long as one finds a way to refine the computation. This is a longstanding challenge in the case of OSp quivers due to a current lack of understanding of such computations.² One of the main results of this paper is the refined Hilbert series, and therefore the enhanced magnetic symmetry of the OSp quivers under our study. Together with the other tools of the Plethystic programme [23, 24], the Hilbert series can be used to give an algebraic description of the moduli space as a variety. We will in particular make use of the notion of highest weight generators (HWGs) developed in [25] in order to write down closed form expressions for the Coulomb branch Hilbert series of the OSp magnetic quivers under our consideration.

In this paper we uncover an interesting phenomenon which is common to all models under our consideration; the moduli space of the OSp quivers that we study generically factorizes into two decoupled sectors, each of which has an alternative description, in terms of the moduli space of a single connected unitary quiver. An upshot of this result is that we can write exact highest weight generating functions (HWGs) encoding the Coulomb

¹One can also consider mixed branches, but we will not explore that in this work.

²This difficulty is related to the notion of hidden FI parameters in OSp quivers, see e.g. [22].

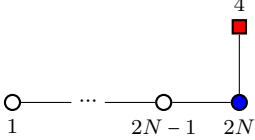
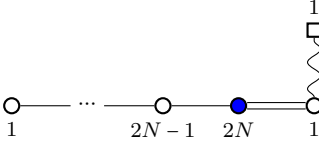
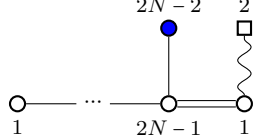
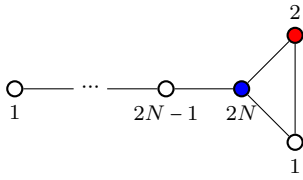
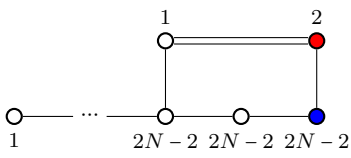
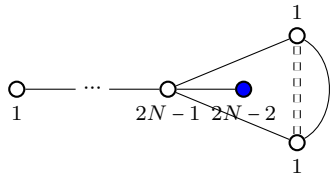
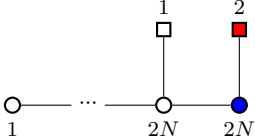
branch Hilbert series of several families of orthosymplectic quivers using known results for the individual factors. This result ultimately follows from fact that the OSp quivers that we study serve as magnetic quivers to 5d $\mathcal{N} = 1$ SCFTs which are the UV fixed point of a 5d IR gauge theory whose gauge group is a product of $SO(4)$ factors, and with matter representations transforming as either spinors or conjugate spinors of each $SO(4)$ factor. Since $SO(4)$ is locally isomorphic to $SU(2) \times SU(2)$, and a spinor and conjugate spinor transform under different $SU(2)$ factors, each such theory can be reformulated as a product of two decoupled theories, each of which has a gauge group that is a product of $SU(2)$ s. The theories containing $SO(4)$ factors can be engineered using a single type IIB brane web with the inclusion of O5-planes, which can then be used to obtain an OSp magnetic quiver [2, 5]. On the other hand the formulation in terms of the product of theories with $SU(2)$ factors is engineered by two independent brane webs, giving rise to two magnetic quivers, which will be unitary by construction [11].

We note that an analogous factorization phenomenon happens for 4d $N=2$ theories of class-S of D-type. In this context, it is well known that there are cases in which a single three-punctured sphere describes the direct sum of 2 SCFTs, each of which also admits a realization in A-type class-S [26–29]. Indeed, we identify some 4d $N=2$ theories of D-type class-S which exhibit this factorization, and for which the orthosymplectic 3d mirror theories correspond to the magnetic quiver derived from the 5-brane webs with O5-plane. Likewise, the 3d mirror of the two A-type factors also corresponds to the unitary magnetic quiver derived from the 5-brane web without O5-plane.

For ease of presentation, we tabulate a list of all the orthosymplectic quivers appearing in this work, along with their Coulomb branch symmetry and the refined highest weight generating functions of the Coulomb branch in table 1.

The organization of this paper is as follows. In section 2 we review the relevant tools that we will need from the plethystic programme. We discuss the monopole formula for the Coulomb branch and the Molien-Weyl formula for the Higgs branch and review the notion of highest weight generating functions. Section 3, contains our main results. Here we will present the product exceptional sequences of OSp quivers, their 5d origin from brane webs with O5 planes and their unitary counterparts as well as their 5d origin from ordinary brane webs. In this section we also state the Hilbert series results for all the quivers. Finally, in section 4 we discuss potential applications of our results and state open problems which we find deserve further investigation. Appendix A contains additional details for some of the Higgs branch Hilbert series computations. We collect the results of HWGs for the unitary quivers in the table 9 of appendix B. In appendix C, we give more details on the computations of the Coulomb branch Hilbert series for the OSp quivers. Appendix D contains a review of the class S technology that is used within the main sections.

Table 1. Summary of the orthosymplectic quivers, their Coulomb branch symmetry and the associated HWG. The corresponding fugacities are denoted by subscripts in the symmetry groups. Note that for $N = 1$, there is an enhancement in the symmetry as detailed in the later sections. The notations for the quivers are defined in the beginning of section 3.

Quiver	Symmetry	PL[HWG]
	$SU(2N)_\mu \times$ $SU(2N)_\nu$	$\sum_{k=1}^N (\mu_k \mu_{2N-k} + \nu_k \nu_{2N-k}) t^{2k}$
	$SU(2N)_\mu \times$ $SU(2N)_\nu \times$ $U(1)_q$	$\sum_{k=1}^N (\mu_k \mu_{2N-k} + \nu_k \nu_{2N-k}) t^{2k} + t^2 +$ $(q + q^{-1}) \nu_N t^{N+1} - \nu_N^2 t^{2N+2}$
	$SU(2N)_\mu \times$ $SU(2N)_\nu \times$ $U(1)_q$	$\sum_{k=1}^N \mu_k \mu_{2N-k} t^{2k} + \sum_{k=1}^{N-1} \nu_k \nu_{2N-k} t^{2k} +$ $t^2 + (\nu_{N+1} q + \nu_{N-1} q^{-1}) t^{N+1} -$ $\nu_{N+1} \nu_{N-1} t^{2N+2}$
	$SU(2N)_\mu \times$ $SU(2N)_\nu \times$ $U(1)_{q_1} \times$ $U(1)_{q_2}$	$\sum_{k=1}^N (\mu_k \mu_{2N-k} + \nu_k \nu_{2N-k}) t^{2k} + 2t^2 +$ $(q_1 + q_1^{-1}) \mu_N t^{N+1} - \mu_N^2 t^{2N+2} +$ $(q_2 + q_2^{-1}) \nu_N t^{N+1} - \nu_N^2 t^{2N+2}$
	$SU(2N)_\mu \times$ $SU(2N)_\nu \times$ $U(1)_{q_1} \times$ $U(1)_{q_2}$	$\sum_{k=1}^{N-1} (\mu_k \mu_{2N-k} + \nu_k \nu_{2N-k}) t^{2k} + 2t^2 +$ $(\mu_{N+1} q_1 + \mu_{N-1} q_1^{-1}) t^{N+1} -$ $\mu_{N+1} \mu_{N-1} t^{2N+2} +$ $(\nu_{N+1} q_2 + \nu_{N-1} q_2^{-1}) t^{N+1} -$ $\nu_{N+1} \nu_{N-1} t^{2N+2}$
	$SU(2N)_\mu \times$ $SU(2N)_\nu \times$ $U(1)_{q_1} \times$ $U(1)_{q_2}$	$\sum_{k=1}^N \mu_k \mu_{2N-k} t^{2k} + \sum_{k=1}^{N-1} \nu_k \nu_{2N-k} t^{2k} +$ $2t^2 + (q_1 + q_1^{-1}) \mu_N t^{N+1} - \mu_N^2 t^{2N+2} +$ $(\nu_{N+1} q_2 + \nu_{N-1} q_2^{-1}) t^{N+1} -$ $\nu_{N+1} \nu_{N-1} t^{2N+2}$
	$SU(2N+1)_\mu \times$ $SU(2N+1)_\nu$	$\sum_{k=1}^N (\mu_k \mu_{2N+1-k} + \nu_k \nu_{2N+1-k}) t^{2k}$

	$\mathrm{SU}(2N+1)_\mu \times \mathrm{SU}(2N+1)_\nu$	$\sum_{k=1}^{N-1} (\mu_k \mu_{2N+1-k} + \nu_k \nu_{2N+1-k}) t^{2k}$
	$\mathrm{SU}(2N+1)_\mu \times \mathrm{SU}(2N+1)_\nu$	$\sum_{k=1}^N \mu_k \mu_{2N+1-k} t^{2k} + \sum_{j=1}^{N-1} \nu_j \nu_{2N+1-j} t^{2j}$
	$\mathrm{SU}(2N+1)_\mu \times \mathrm{SU}(2N+1)_\rho \times \mathrm{SU}(2)_\nu \times \mathrm{U}(1)_q$	$\sum_{i=1}^N (\mu_i \mu_{2N+1-i} + \rho_i \rho_{2N+1-i}) t^{2i} + (\nu^2 + 1) t^2 + \nu (\mu_N q + \mu_{N+1} q^{-1}) t^{N+1} - \nu^2 \mu_N \mu_{N+1} t^{2N+2}$
	$\mathrm{SU}(2N+1)_\mu \times \mathrm{SU}(2N+1)_\lambda \times \mathrm{SU}(2)_\nu \times \mathrm{SU}(2)_\eta \times \mathrm{U}(1)_q$	$\sum_{i=1}^N (\mu_i \mu_{2N+1-i} + \lambda_i \lambda_{2N+1-i}) t^{2i} + (\nu^2 + 1) t^2 + \nu (\mu_N q + \mu_{N+1} q^{-1}) t^{N+1} - \nu^2 \mu_N \mu_{N+1} t^{2N+2} + \eta^2 t^2$
	$\mathrm{SU}(2N+1)_\mu \times \mathrm{SU}(2N+1)_\eta \times \mathrm{SU}(2)_\nu \times \mathrm{SU}(2)_\lambda \times \mathrm{U}(1)_q \times \mathrm{U}(1)_r$	$\sum_{i=1}^N (\mu_i \mu_{2N+1-i} + \eta_i \eta_{2N+1-i}) t^{2i} + (\nu^2 + \lambda^2 + 2) t^2 + \nu (\mu_N q + \mu_{N+1} q^{-1}) t^{N+1} + \lambda (\eta_N r + \eta_{N+1} r^{-1}) t^{N+1} - \nu^2 \mu_N \mu_{N+1} t^{2N+2} - \lambda^2 \eta_N \eta_{N+1} t^{2N+2}$

	$\begin{aligned} & \text{SU}(2N+2)_\mu \times \\ & \text{SU}(2N+2)_\lambda \times \\ & \text{SU}(2)_{\nu_1} \times \\ & \text{SU}(2)_{\nu_2} \times \\ & \text{SU}(2)_{\rho_1} \times \\ & \text{SU}(2)_{\rho_2} \end{aligned}$	$\sum_{i=1}^{N+1} (\mu_i \mu_{2N+2-i} + \lambda_i \lambda_{2N+2-i}) t^{2i} + (\nu_1^2 + \nu_2^2 + \rho_1^2 + \rho_2^2) t^2 + 2t^4 + \nu_1 \nu_2 \mu_{N+1} (t^{N+1} + t^{N+3}) + \rho_1 \rho_2 \lambda_{N+1} (t^{N+1} + t^{N+3}) - \nu_1^2 \nu_2^2 \mu_{N+1}^2 t^{2N+6} - \rho_1^2 \rho_2^2 \lambda_{N+1}^2 t^{2N+6}$
	$\begin{aligned} & \text{SO}(4N+6)_\mu \times \\ & \text{SO}(4N+6)_\nu \end{aligned}$	$\sum_{k=1}^N (\mu_{2k} + \nu_{2k}) t^{2k}$
	$\begin{aligned} & \text{SO}(4N+6)_\mu \times \\ & \text{SO}(4N+6)_\nu \times \\ & \text{U}(1)_q \end{aligned}$	$\sum_{i=1}^N (\mu_{2i} + \nu_{2i}) t^{2i} + t^2 + (\mu_{2N+2q} + \mu_{2N+3q} q^{-1}) t^{N+1}$

	$\text{SO}(4N+6)_\mu \times \text{SO}(4N+6)_\nu \times \text{U}(1)_{q_1} \times \text{U}(1)_{q_2}$	$\sum_{i=1}^N (\mu_{2i} + \nu_{2i}) t^{2i} + 2t^2 + (\mu_{2N+2} q_1 + \mu_{2N+3} q_1^{-1}) t^{N+1} + (\nu_{2N+2} q_2 + \nu_{2N+3} q_2^{-1}) t^{N+1}$
	$\text{SO}(4N+8)_\mu \times \text{SO}(4N+8)_\lambda \times \text{SU}(2)_\nu \times \text{SU}(2)_\rho$	$\sum_{i=1}^{N+1} (\mu_{2i} + \lambda_{2i}) t^{2i} + 2t^4 + (\nu^2 + \rho^2) t^2 + \nu \mu_{2N+4} (t^{N+1} + t^{N+3}) + \rho \lambda_{2N+4} (t^{N+1} + t^{N+3}) + (\mu_{2N+4}^2 + \lambda_{2N+4}^2) t^{2N+4} - \nu^2 \mu_{2N+4} t^{2N+6} - \rho^2 \lambda_{2N+4} t^{2N+6}$

2 Preliminaries: tools from the plethystic programme

In this section we review the material we need for the computation of the Hilbert series. The discussion will be minimal and will cover only those aspects necessary for the subsequent section. For more details the reader can consult the original papers. The literature on this material is vast, but we will mostly follow [6, 20, 21, 23, 25]. In the following subsections we consider a 3d $\mathcal{N} = 4$ theory with gauge group G of rank r and n_H hypermultiplets transforming under the representation R_i of G ($i = 1, \dots, n_H$).

2.1 Coulomb branches

The bosonic fields in a 3d $\mathcal{N} = 4$ vector multiplet consist of a gauge field and 3 real scalars. Upon dualising the gauge field to a scalar we have 4 scalars at our disposal, which can be pairwise complexified. These are the coordinates on the Coulomb branch at large vevs. Each of these complex scalars forms the scalar component of an $\mathcal{N} = 2$ chiral superfield. For small vevs one needs to replace the complex scalar containing the dual photon by a BPS monopole operator [16], which can again be thought of as the bottom component of an $\mathcal{N} = 2$ chiral multiplet. The quantum coordinates of the Coulomb branch are therefore the BPS monopole V_m of magnetic charge m and the adjoint scalar ϕ . The Hilbert series (HS) for the Coulomb branch of a good or ugly (in the sense of [16]) 3d $\mathcal{N} = 4$ theory is

then computed by the monopole formula [20]

$$\text{HS}(t, z) = \sum_{m \in \Lambda_m^{G^\vee} / \mathcal{W}_{G^\vee}} z^{J(m)} t^{2\Delta(m)} P_G(t, m), \quad (2.1)$$

where the sum is over the magnetic lattice of the gauge group G , we refer to [6] for a detailed account. $\Delta(m)$ is the conformal dimension of the monopole operator with magnetic charge m , and is given by

$$\Delta(m) = - \sum_{\alpha \in \Delta^+} |\alpha(m)| + \frac{1}{2} \sum_{i=1}^{n_H} \sum_{\rho_i \in R_i} |\rho_i(m)|, \quad (2.2)$$

where Δ^+ is the set of positive roots and ρ_i are the weights of the representation R_i . The classical dressing factor $P_G(t, m)$ counts invariants built out of the adjoint scalars ϕ

$$P_G(t, m) = \prod_{i=1}^r \frac{1}{1 - t^{2d(i)}}, \quad (2.3)$$

where $d(i)$ are the degrees of the Casimir invariants of the gauge symmetry $H \subset G$ which is the unbroken part of the original gauge symmetry in the presence of a monopole operator V_m of charge m . Finally z denote the fugacities of the topological symmetry whose exponent $J(m)$ denotes the topological current.

2.2 Higgs branches

The Higgs branch of a 3d $\mathcal{N} = 4$ theory is parametrised by vevs of the scalar components of the hypermultiplets. When dealing with a gauge theory, each such operator will be in an irrep of G . The Higgs branch operators are therefore those constructed by considering all symmetrised tensor powers of these irreps. The symmetrisation is done in order to be consistent with Pauli statistics. To avoid overcounting, the relations among these scalar operators due to the superpotential need to be imposed. One will then need to project onto the gauge singlet states in order for the resulting operators to be well defined gauge invariant operators. The Higgs branch Hilbert series is therefore computed using the Molien-Weyl formula [21]

$$\text{HS}^{\mathcal{H}}(t) = \int_G d\mu_G \frac{\text{PE} \left[\sum_{i=1}^{n_H} \chi_{R_i}(x) t \right]}{\text{PE} \left[\chi_{\mathbf{Adj}}(x) t^2 \right]}, \quad (2.4)$$

where $\chi_{R_i}(x)$ is the character of the representation R_i of G under which the scalars in the i -th hypermultiplet transform, $\chi_{\mathbf{Adj}}$ is the character of the adjoint representation of G , the representation carried by the relations. The function $\text{PE}[\cdot]$ is the plethystic exponential, defined via

$$\text{PE} [f(z_1, \dots, z_r)] = \exp \left(\sum_{k=1}^{\infty} \frac{1}{k} f(z_1^k, \dots, z_r^k) \right), \quad (2.5)$$

it is a symmetrising function that generates the characters of the symmetrised tensor powers of χ_{R_i} . Finally the projection onto gauge invariant operators is achieved by integrating over the group manifold using the Haar measure. In a suitable basis, the Haar measure can be taken to be

$$\int_G d\mu_G = \frac{1}{(2\pi i)^r} \prod_{i=1}^r \oint_{|x_i|=1} \frac{dx_i}{x_i} \prod_{\alpha \in \Delta^+} \left(1 - \prod_{k=1}^r x_k^{\alpha_k} \right). \quad (2.6)$$

2.3 Highest weight generating functions

The refined Hilbert series for a given theory can be generically expanded as Taylor series in t such that the coefficients are sum of the characters of representations of global symmetry of the theory. In general, given a global symmetry of rank r , the refined Hilbert series can be expanded as:

$$\sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} \cdots \sum_{n_r=0}^{\infty} \chi_{[f_1, f_2, \dots, f_r]} t^{\eta}, \quad (2.7)$$

where each of f_1, \dots, f_r and η can be some polynomial function in variables n_1, \dots, n_r . The notation $\chi_{[f_1, \dots, f_r]}$ is the character of the irrep of the global symmetry whose highest weight is $f_1 \Lambda_1 + \dots + f_r \Lambda_r$, where f_i are the Dynkin labels of the irrep and Λ_i are the fundamental weights of the global symmetry group. A convenient way to repackage the same information is in terms of highest weight generating functions or HWGs [25]. One introduces a set of fugacities $\{\mu_1, \dots, \mu_r\}$, called highest weight fugacities, and one writes the characters in terms of μ_i according to the map

$$\chi_{[f_1, \dots, f_r]} \leftrightarrow \mu_1^{f_1} \cdots \mu_r^{f_r}. \quad (2.8)$$

With this map, the Hilbert series becomes a formal power series which can be resummed, the corresponding generating function is termed as its HWG:

$$\text{HWG} = \sum_{n_1=0}^{\infty} \cdots \sum_{n_r=0}^{\infty} \mu_1^{f_1} \cdots \mu_r^{f_r} t^{\eta}. \quad (2.9)$$

3 Product sequences

In this section we present sequences of orthosymplectic magnetic quivers whose moduli space is the product of two decoupled sectors, each of which enjoys a description as the moduli space of a unitary quiver. Each sequence is parameterised by an integer N which is the sequence number, and labelled $E_m \times E_n$, in accordance with the Coulomb branch isometry of the $N = 1$ case. For $N = 1$, some of the theories become particularly simple such that we can prove the equivalence of the orthosymplectic quiver with the two unitary quivers. The $E_m \times E_n$ orthosymplectic quiver with sequence number 1 corresponds to the magnetic quiver for infinite coupling limit of 5d $\mathcal{N} = 1$ SO(4) gauge theory with $m - 1$ hypermultiplets in the spinor representation denoted by \mathbf{s} and $n - 1$ hypermultiplets in the conjugate spinor representation of SO(4) denoted by \mathbf{c} . Correspondingly the dual unitary quivers with sequence number 1 correspond to magnetic quivers for infinite coupling limit of 5d $\mathcal{N} = 1$ SU(2)×SU(2) gauge theory with $m - 1$ hypermultiplets in the $(\mathbf{F}, \mathbf{1})$ and $n - 1$ hypermultiplets in the $(\mathbf{1}, \mathbf{F})$ representation of SU(2)×SU(2), where we denote by \mathbf{F} the fundamental representation of associated gauge group. One can engineer these theories using 5-brane webs with O5-planes [30], as well as using ordinary brane webs [31]. This pattern generalizes for higher sequence numbers, namely one can provide an intuition for the reason that the orthosymplectic quivers factorise into two decoupled sectors by viewing them as magnetic quivers of a 5d theory. We will therefore employ this perspective in the following. We will use EQ $_{m,n}$ to denote the 5d OSp electric quivers for the $E_m \times E_n$

sequence, while we use the notation EQ_m to denote the 5d unitary electric quivers to which the former factorise. Similarly, we will denote the OSp magnetic quivers of the $E_m \times E_n$ sequence by $\text{MQ}_{m,n}$, while we denote the unitary components to which they factorise by MQ_m . Occasionally there will be more than one generalisation of a given sequence for higher sequence numbers, in which case we will distinguish the different sequences by a prime.

Our conventions for the magnetic quivers are identical to those appearing in [2] which we briefly review in the following. We denote by a white, red and blue node, the groups $\text{U}(n)$, $\text{SO}(n)$ and $\text{USp}(2n)$ respectively. A circular node is to be understood as a gauge group, while a square node denotes a flavour group. We use solid lines to denote bifundamental hypermultiplets, in the case where the solid line connects an orthogonal to a symplectic node there is a reality condition which renders the link to be a half hypermultiplet. We use a dashed link between two unitary nodes to denote a fundamental-fundamental hypermultiplet and a jagged line between a unitary flavour and a gauge $\text{U}(1)$ node to denote a charge 2 hypermultiplet.

3.1 The $E_1 \times E_1$ sequence

Consider the 5-brane web constructed by collapsing $2N$ NS5 branes on top of an O5-plane that is asymptotically O5^+ as in figure 1. By resolving this web to go on the Coulomb branch one can identify the following low energy quiver description,

$$\text{EQ}_{1,1} = \overbrace{\text{SO}(4) - \text{USp}(0) - \text{SO}(4) - \dots - \text{USp}(0) - \text{SO}(4)}^{2N-1} . \quad (3.1)$$

One can recast this theory, by using the accidental Lie algebra isomorphism $\mathfrak{so}(4) \cong \mathfrak{su}(2) \oplus \mathfrak{su}(2)$ as a direct sum of two decoupled 5d quiver gauge theories. For each $\text{SU}(2)$ gauge node, we need to fix the discrete theta angles, $\theta = 0$ or $\theta = \pi$. It is known that the former has a Higgs branch with dimension one while the latter does not have Higgs branch in the infinite coupling limit. Since the $\text{SO}(4)$ gauge theory constructed as in the left of figure 1 with $N = 1$ has Higgs branch with dimension two in the infinite coupling limit, it should correspond to $\text{SU}(2)_0 \times \text{SU}(2)_0$ gauge theory. Further support for this claim is given in appendix E by constructing the $\text{SU}(2)_{\theta_1} \times \text{SU}(2)_{\theta_2}$ gauge theories with different discrete theta angles by using 5-brane web diagram with O5-planes. Since each $\text{SO}(4)$ gauge node corresponds to $\text{SU}(2)_0 \times \text{SU}(2)_0$, the generalization to generic N should be³

$$\text{EQ}_{1,1} = (\text{EQ}_1)^2 = \left(\overbrace{\text{SU}(2)_0 - \text{SU}(2)_0 - \text{SU}(2)_0 - \dots - \text{SU}(2)_0}^N \right)^2 . \quad (3.2)$$

Since this description involves only special unitary gauge groups, we should be able to engineer it using (two copies of) ordinary brane webs, i.e. without using O5-planes. It is

³Analogous discussion on the discrete theta angle will apply to the later examples. In the following subsections, we suppress the label “0” of $\text{SU}(2)_0$ denoting its discrete theta angle.

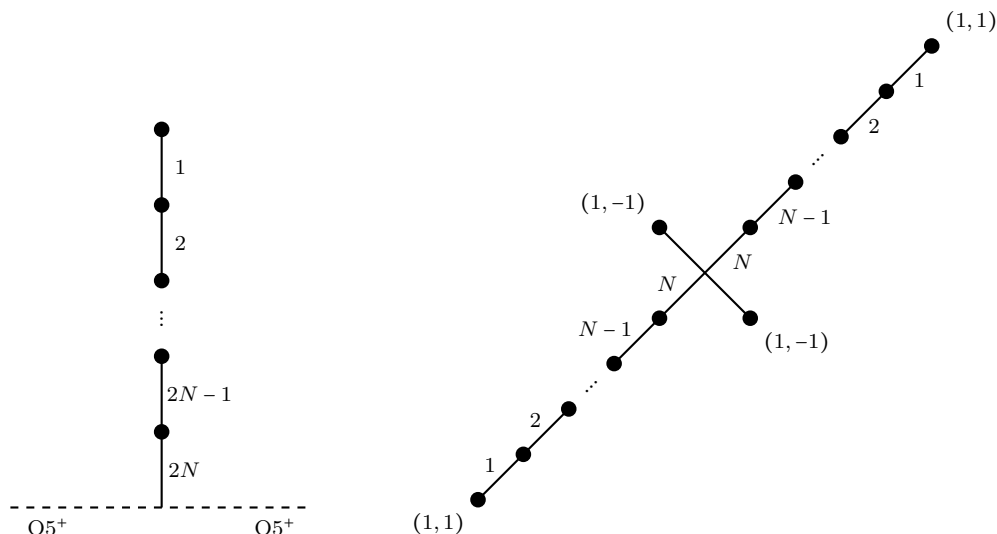


Figure 1. Brane webs engineering $\text{EQ}_{1,1}$ (left) and EQ_1 (right). The numbers near each 5-brane denote the number of coincident 5-branes in the stack in that segment. Black dots denote 7-branes of charge (p, q) .

constructed in such a way that it is decomposed into the copy of $SU(2)_0$ gauge theories in the region where the bi-fundamental masses are large enough [32]. The result is given in the right of figure 1. This is also discussed in appendix E.

At this point we need to clarify which aspect of the two theories (3.1) and (3.2) are expected to be the same. This is because we used an isomorphism at the level of Lie algebra, ignoring any issues related to the global structure of the gauge group (with the exception of the choice of discrete theta angle mentioned above). In particular, any information related to local operators in the two theories is likely to agree, while questions about e.g. line and surface operators in general will be sensitive to the global structure of the gauge group. Our primary interest in these theories is in their Higgs branch, which is parameterised by local operators, and thus should agree regardless of any subtle differences such as those alluded to above.

Having constructed web diagrams for $\text{EQ}_{1,1}$ and EQ_1 in figure 1, we can now proceed to derive their magnetic quivers, following the rules introduced in [2, 5, 11]. From the orientifold web in figure 1 we obtain the OSp magnetic quiver⁴

$$\text{MQ}_{1,1} = \begin{array}{c} & & 4 \\ & & \blacksquare \\ & | \\ \circ - \dots - \circ - \bullet \\ 1 \qquad\qquad 2N-1 \quad 2N \end{array}, \quad (3.3)$$

⁴The flavor $\text{SO}(4)$ node appearing in the quiver $\text{MQ}_{1,1}$ was argued [2] to arise from the intersection between $(0,1)$ 5-brane and $\text{O}5^+$ -plane. This is based on the intuition that an $\text{O}5^+$ carries the same charge as an $\text{O}5^-$ and two (immobile) D5 branes.

while the characters for the fundamental and second rank antisymmetric representation of $\mathrm{USp}(4)$ are given respectively by⁵

$$\begin{aligned}\chi_{[1,0]_{C_2}} &= x_1 + \frac{x_2}{x_1} + \frac{x_1}{x_2} + \frac{1}{x_1}, \\ \chi_{[0,1]_{C_2}} &= x_2 + \frac{x_1^2}{x_2} + 1 + \frac{x_2}{x_1^2} + \frac{1}{x_2},\end{aligned}\tag{3.9}$$

Thus the expression we need to evaluate is

$$\begin{aligned}& \prod_{q=2}^4 (1 - t^{2q}) \oint_{|x_1|=1} \frac{dx_1}{2\pi i x_1} \oint_{|x_2|=1} \frac{dx_2}{2\pi i x_2} (1 - x_1^2)(1 - x_2) \left(1 - \frac{x_1^2}{x_2}\right) \left(1 - \frac{x_2^2}{x_1^2}\right) \\ & \times \frac{1}{(1 - x_2 t^2) \left(1 - \frac{x_1^2}{x_2} t^2\right) (1 - t^2) \left(1 - \frac{x_2^2}{x_1^2} t^2\right) \left(1 - \frac{t^2}{x_2}\right) (1 - x_1 t)^4 \left(1 - \frac{x_2}{x_1} t\right)^4 \left(1 - \frac{x_1}{x_2} t\right)^4 \left(1 - \frac{t}{x_1}\right)^4}.\end{aligned}\tag{3.10}$$

This integral can now be evaluated using residues to arrive at the following Hilbert series

$$\mathrm{HS}_{1,1}^{\mathcal{H}}|_{N=2} = \frac{(1 - t^6)^2 (1 - t^8)^2}{(1 - t^4)^6 (1 - t^2)^6}.\tag{3.11}$$

We recognise this as the Coulomb branch Hilbert series of two copies of $\mathrm{U}(2)$ with 4 fundamental hypermultiplets [20], which is the mirror of the $N = 2$ quiver of (3.4). Therefore we see that the agreement between the OSp quiver (3.3) and unitary quiver (3.4) holds also on the Higgs branch.

3.2 The $E_1 \times E_3$ sequence

In the previous subsection, we saw that the Higgs branch of the fixed point limit of $5d$ $\mathrm{SO}(4)$ gauge theory, factorises to two copies of the Higgs branch of the $5d$ pure $\mathrm{SU}(2)$ gauge theory. It is natural to ask whether this pattern holds if we include matter transforming under $\mathrm{SO}(4)$. The two matter representations which are straightforward to obtain from the brane web are the vector of $\mathrm{SO}(4)$ and the two spinor representations of opposite chirality. Since the vector of $\mathrm{SO}(4)$ corresponds to bifundamental of $\mathrm{SU}(2) \times \mathrm{SU}(2)$, this will not lead to the desired factorised theory. However, each of the two spinor representations, denoted by \mathbf{s} and \mathbf{c} respectively, will only transform under one of the two $\mathrm{SU}(2)$ factors in $\mathrm{SO}(4)$. In this and subsequent subsections, we will exploit this well-known fact.

Consider the orientifold web diagram presented in figure 2. The corresponding IR quiver gauge theory description is given by the electric quiver

$$\mathrm{EQ}_{1,3} = [1 \ \mathbf{s}] - \overbrace{\mathrm{SO}(4) - \mathrm{USp}(0) - \mathrm{SO}(4) - \dots - \mathrm{USp}(0) - \mathrm{SO}(4)}^{2N-1} - [1 \ \mathbf{s}].\tag{3.12}$$

By using the isomorphism $\mathfrak{so}(4) \cong \mathfrak{su}(2) \times \mathfrak{su}(2)$, we can rewrite this theory as a product

⁵These characters are computed as follows. For a weight $w = [w_1, w_2] \equiv w_1 \Lambda_1 + w_2 \Lambda_2$ appearing in the weight system of a representation R of C_2 , the corresponding monomial in the character will be $x_1^{w_1} x_2^{w_2}$. For example, the weights appearing in the weight system of fundamental representation of C_2 are $\{[1, 0], [-1, 1], [1, -1], [-1, 0]\}$ where each weight is written in the fundamental weight basis. Thus the character for fundamental representation will be simply $x_1^1 x_2^0 + x_1^{-1} x_2^1 + x_1^1 x_2^{-1} + x_1^{-1} x_2^0$.

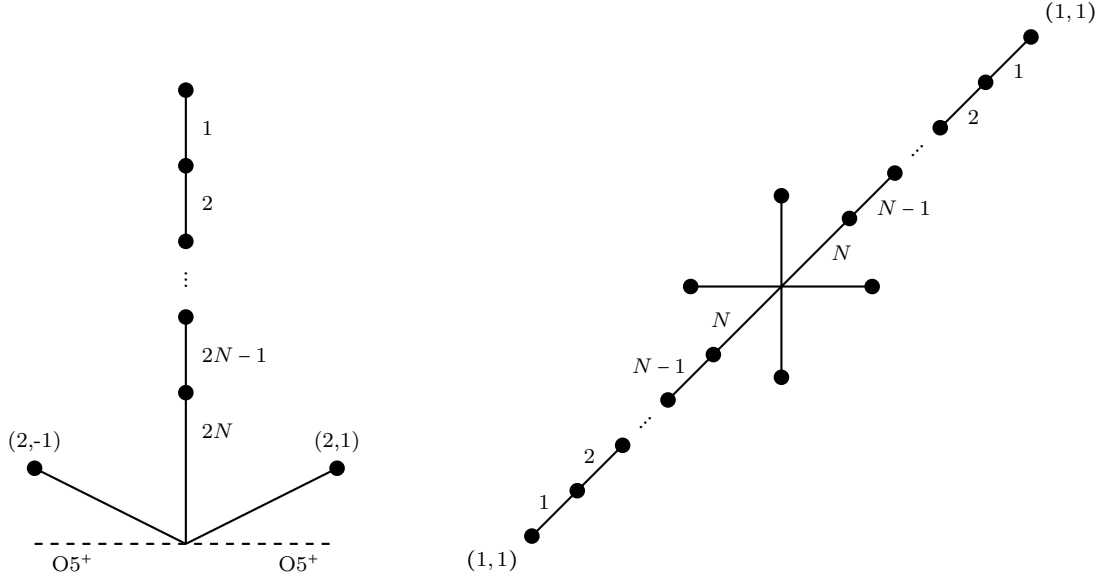


Figure 2. Brane webs engineering $EQ_{1,3}$ (left) and EQ_3 (right). For the unitary brane web engineering EQ_1 , see figure 1.

of the following electric quivers.

$$EQ_{1,3} = EQ_1 \times EQ_3 = \overbrace{SU(2) - SU(2) - \dots - SU(2)}^N \times \overbrace{\begin{array}{c} SU(2) - SU(2) - \dots - SU(2) \\ | \qquad \qquad \qquad | \\ [1\mathbf{F}] \qquad \qquad \qquad [1\mathbf{F}] \end{array}}^N \quad (3.13)$$

The unitary brane web for EQ_1 is given in figure 1, while the unitary brane web for EQ_3 is presented in figure 2. The orientifold web in figure 2 admits two maximal subdivisions. Accordingly the Higgs branch of this theory is the union of two cones, given by the two OSp magnetic quivers in table 2. On the other hand, we expect these magnetic quivers to be equivalent to the product $MQ_1 \times (MQ_3^{(I)} \cup MQ_3^{(II)})$, with the latter factor obtained from the unitary web for EQ_3 in figure 2. As a further non-trivial check, we can compute the Coulomb branch Hilbert series of the unitary and OSp quivers. The unitary quivers MQ_1 along with $MQ_3^{(I)}$ and $MQ_3^{(II)}$ have known HWGs [3], which we have tabulated in table 9.

We are now in a position to write down the HWGs for the OSp magnetic quivers in table 2. The final result for the first cone reads

$$HWG_{1,3}^{(I)} = PE \left[\sum_{k=1}^N \mu_k \mu_{2N-k} t^{2k} \right] PE \left[t^2 + (q + q^{-1}) \nu_N t^{N+1} + \sum_{k=1}^N \nu_k \nu_{2N-k} t^{2k} - \nu_N^2 t^{2N+2} \right], \quad (3.14)$$

where μ and ν are the highest weight fugacities for $SU(2N) \times SU(2N)$ while q keeps track

MS	OSp	Unitary
$MQ_{1,3}^{(I)}$		
$MQ_{1,3}^{(II)}$		

Table 2. OSp and unitary representation of the two cones on the Higgs branch of $EQ_{1,3}$. The unitary quivers appearing in the extreme right of the two rows of the table are respectively $MQ_3^{(I)}$ and $MQ_3^{(II)}$.

of the $U(1)$ charge. The HWG for the second cone reads

$$\begin{aligned}
 \text{HWG}_{1,3}^{(II)} = & \text{PE} \left[\sum_{k=1}^N \mu_k \mu_{2N-k} t^{2k} \right] \\
 & \times \text{PE} \left[t^2 + (\nu_{N+1} q + \nu_{N-1} q^{-1}) t^{N+1} + \sum_{k=1}^{N-1} \nu_k \nu_{2N-k} t^{2k} - \nu_{N+1} \nu_{N-1} t^{2N+2} \right]. \quad (3.15)
 \end{aligned}$$

This can be verified upon comparison with the result of an unrefined Hilbert series computation on the OSp side. The results for low values of N are given in table 10.

We can also compute the Higgs branch HS for $MQ_{1,3}^{(I)}$ for $N = 2$ exactly. The computation is very similar to that of the Higgs branch of $MQ_{1,1}$ discussed around (A.2). We need to evaluate the following integral

$$\prod_{q=1}^4 (1 - t^{2q}) \int d\mu_{C_2} \int d\mu_{U(1)} \text{PE} \left[\chi_{[0,1]}^{C_2} t^2 + 2\chi_{[1,0]}^{C_2} (q + q^{-1}) t + (q^2 + q^{-2}) t \right] \quad (3.16)$$

Evaluating this integral by finding the residues one arrives at the following

$$\text{HS}_{1,3,(I)}^{\mathcal{H}} = \frac{(1 - t + t^2)(1 + t^4)(1 + t^3 + t^4 + t^5 + t^6 + t^7 + t^{10})}{(1 - t)^8 (1 + t)^6 (1 + t^2)^3 (1 + t + t^2 + t^3 + t^4)}. \quad (3.17)$$

This is to be compared with the product of the Higgs branch HS of the two unitary quivers appearing in the first row of table 2. We already know the result for one of these, which is identical to MQ_1 of (3.4). Its Higgs branch HS was discussed in the previous section and is given by the square root of the expression in (3.11). The Higgs branch HS of the

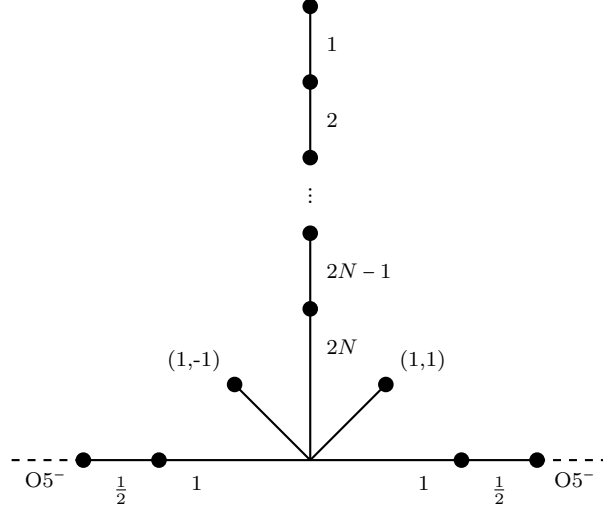


Figure 3. Orientifold web for $\text{EQ}_{3,3}$.

other quiver in the first row of table, 2, which we dub $\text{MQ}_3^{(1)}$ is straightforward to compute. Specialising to the case $N = 2$, we need to evaluate the following integral

$$\begin{aligned}
 & \int d\mu_{\text{U}(2)}(x, q_x) \int d\mu_{\text{U}(1)}(u) H_{T[\text{SU}(2)]}^2(x) H_{\text{glue}}^{\text{U}(2)}(x, q_x) \\
 & \times H_{[2]-[1]}(x, q_x) H_{[2]-[1]}(x, q_x, u) H_{[1]-[1]}(u) H_{\text{glue}}^{\text{U}(1)}(u) \\
 & = \oint_{|x|=1} \frac{dx}{2\pi i x} (1-x^2) \oint_{|q_x|=1} \frac{dq_x}{2\pi i q_x} \oint_{|u|=1} \frac{du}{2\pi i u} (1-t^2)^2 (1-t^4)^2 \\
 & \times \text{PE} \left[(x^2 + 1 + x^{-2})t^2 + (x + x^{-1})(q_x + q_x^{-1})(u + u^{-1})t + (u + u^{-1})t + (x + x^{-1})(q_x + q_x^{-1})t \right].
 \end{aligned} \tag{3.18}$$

Evaluating this integral by computing its residues results in

$$\text{HS}_{3,(1)}^{\mathcal{H}} = \frac{(1 + t^3 + t^4 + t^5 + t^6 + t^7 + t^{10})}{(1-t)^4 (1+t)^2 (1+t^2) (1+t+t^2) (1+t+t^2+t^3+t^4)}. \tag{3.19}$$

Together with the result for the HB of MQ_1 , this precisely reproduces the computation on the OSp side (3.17).

3.3 The $E_3 \times E_3$ sequence

The $E_3 \times E_3$ sequence corresponds to the fixed point limit of the electric quiver

$$\text{EQ}_{3,3} = (\text{EQ}_3)^2 = \overbrace{\text{SO}(4) - \text{USp}(0) - \text{SO}(4) - \dots - \text{USp}(0) - \text{SO}(4)}^{2N-1} \begin{matrix} | \\ [1\mathbf{s} + 1\mathbf{c}] \end{matrix} \begin{matrix} | \\ [1\mathbf{s} + 1\mathbf{c}] \end{matrix}. \tag{3.20}$$

The orientifold web which engineers this theory is given in figure 3. This brane web admits three maximal subdivisions leading to the three OSp magnetic quivers in table 3. It can be

MS	OSp	Unitary
$\text{MQ}_{3,3}^{(I)}$		$\left(\begin{array}{c} \text{Unitary web with } N \text{ nodes} \end{array} \right)^2$
$\text{MQ}_{3,3}^{(II)}$		$\left(\begin{array}{c} \text{Unitary web with } N-1 \text{ nodes} \end{array} \right)^2$
$\text{MQ}_{3,3}^{(III)}$		

Table 3. Unitary and OSp magnetic quivers for the $E_3 \times E_3$ sequence.

understood as a limiting case of the $Y_N^{1,1}$ theory in [2]. One can provide a purely unitary description of this theory in terms of the following electric quiver:

$$\text{EQ}_{3,3} = \text{EQ}_3^2 = \left(\begin{array}{c} \overbrace{\text{SU}(2) - \text{SU}(2) - \dots - \text{SU}(2)}^N \\ \text{[1F]} \quad \quad \quad \text{[1F]} \end{array} \right)^2, \quad (3.21)$$

which can be engineered by taking two copies of the unitary web shown in figure 2. This brane web admits two maximal subdivisions whose magnetic quivers were discussed in the previous subsection. Since we are taking two copies, a third cone arises when we take a different maximal subdivision for each web diagram. This leads us to the unitary magnetic quivers in table 3.

We can now infer the HWG for the OSp quivers in table 3 by taking those of the corresponding unitary magnetic quivers as building blocks. This reasoning leads us to conjecture the following HWG for the three cones in table 3

$$\begin{aligned} \text{HWG}_{3,3}^{(I)}(t^2) &= \text{PE} \left[t^2 + (\mu_N q_1 + \mu_N q_1^{-1}) t^{N+1} + \sum_{k=1}^N (\mu_k \mu_{2N-k} t^{2k}) - \mu_N^2 t^{2N+2} \right] \\ &\times \text{PE} \left[t^2 + (\nu_N q_2 + \nu_N q_2^{-1}) t^{N+1} + \sum_{k=1}^N (\nu_k \nu_{2N-k} t^{2k}) - \nu_N^2 t^{2N+2} \right] \end{aligned}$$

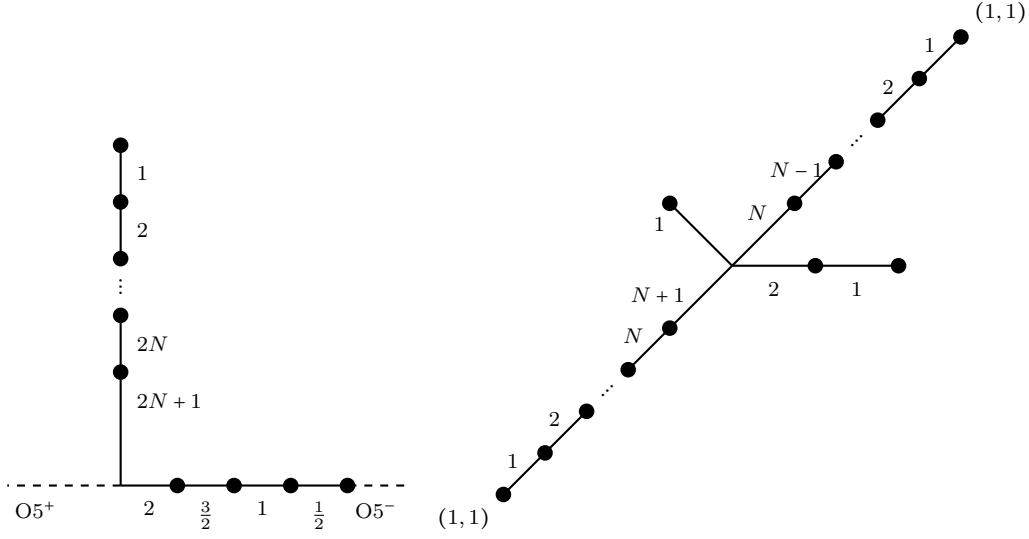


Figure 4. Brane webs for $\text{EQ}_{3',3'}$ (left) and $\text{EQ}_{3'}$ (right).

$$\begin{aligned}
 \text{HWG}_{3,3}^{(\text{II})}(t^2) &= \text{PE} \left[t^2 + (\mu_{N+1}p + \mu_{N-1}p^{-1})t^{N+1} + \sum_{k=1}^{N-1} \mu_k \mu_{2N-k} t^{2k} - \mu_{N+1} \mu_{N-1} t^{2N+2} \right] \\
 &\quad \times \text{PE} \left[t^2 + (\nu_{N+1}q + \nu_{N-1}q^{-1})t^{N+1} + \sum_{k=1}^{N-1} \nu_k \nu_{2N-k} t^{2k} - \nu_{N+1} \nu_{N-1} t^{2N+2} \right] \\
 \text{HWG}_{3,3}^{(\text{III})}(t^2) &= \text{PE} \left[t^2 + (\mu_{N+1}p + \mu_{N-1}p^{-1})t^{N+1} + \sum_{k=1}^{N-1} \mu_k \mu_{2N-k} t^{2k} - \mu_{N+1} \mu_{N-1} t^{2N+2} \right] \\
 &\quad \times \text{PE} \left[t^2 + (q + q^{-1})\nu_N t^{N+1} + \sum_{k=1}^N \nu_k \nu_{2N-k} t^{2k} - \nu_N^2 t^{2N+2} \right]. \quad (3.22)
 \end{aligned}$$

Here μ and ν are the fugacities for the two $\text{SU}(2N)$ groups and p and q are the $\text{U}(1)$ charges. We have verified this result by an explicit unrefined Hilbert series computation of the Coulomb branch of the OSp quivers which are presented in table 10 for low values of N .

3.4 The $E'_3 \times E'_3$ sequence

There is another sequence whose first member has an $E_3 \times E_3$ symmetry. We will refer to this as the $E'_3 \times E'_3$ sequence. In figure 4 we present the orientifold web that engineers the 5d electric quivers in this sequence, the IR quiver description is given by

$$\text{EQ}_{3',3'} = \overbrace{\text{SO}(4) - \text{USp}(0) - \text{SO}(4) - \cdots - \text{USp}(0) - \text{SO}(4)}^{2N-1} \quad (3.23)$$

$\begin{array}{c} | \\ [2s+2c] \end{array}$

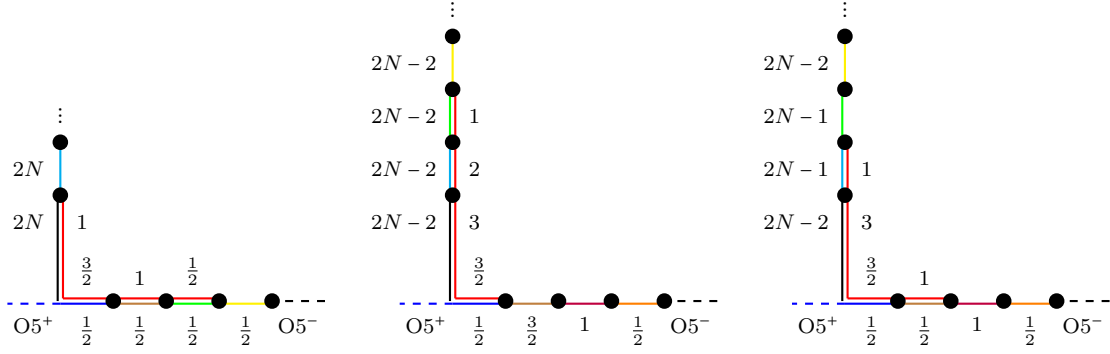


Figure 5. Maximal subdivisions of the Higgs branch of $\text{EQ}_{3',3'}$ at the superconformal limit. The subweb coloured in red is frozen and contributes only as flavour nodes to the magnetic quivers in table 4.

It can also be given a unitary web description, once rewritten as a quiver theory with $\text{SU}(2)$ nodes:

$$\text{EQ}_{3',3'} = \text{EQ}_{3'}^2 = \left(\overbrace{\text{SU}(2) - \text{SU}(2) - \dots - \text{SU}(2)}^N \right)^2_{[2\mathbf{F}]} . \quad (3.24)$$

We present the unitary web engineering each EQ_3' factor in figure 4. Now we can use the unitary and orientifold webs in figure 4 to obtain unitary and OSp magnetic quivers for the $E_3' \times E_3'$ sequence, which appear in table 4. These MQs can be obtained by considering the maximal subdivisions appearing in figure 5. The first subdivision was already discussed in [2], while the second and third were overlooked. It was noticed in [9] that there should be 2 additional OSp cones to match the analysis on the unitary side. We claim that the missing cones in this case correspond to the two new maximal subdivisions appearing in figure 5.

Notice that for $N = 1$, the relation between the unitary and OSp quivers in the first row of table 4 was already suggested in [33], our result generalises this to higher N . The HWG for the unitary quiver is known and appears in [3]. Given the correspondence between the unitary and OSp magnetic quivers in table 4, we can use the results for the HWGs of the unitary quivers to obtain the HWGs for the OSp quivers. In order to do this, let us point out the following useful fact; one of the unitary quivers appearing in the second and third row of table 5, is itself a product of two unitary quivers:

$$\begin{array}{c} \text{Quiver 1} \end{array} = \begin{array}{c} \text{Quiver 2} \end{array} \times \begin{array}{c} \text{Quiver 3} \end{array}, \quad (3.25)$$

MS	OSp	Unitary
$\text{MQ}_{3',3'}^{(\text{I})}$		
$\text{MQ}_{3',3'}^{(\text{II})}$		
$\text{MQ}_{3',3'}^{(\text{III})}$		

Table 4. Magnetic quivers for the $E'_3 \times E'_3$ sequence.

where the right hand side of the above is obtained after ungauging the overall decoupled $U(1)$ in the original quiver. The first quiver in the right hand side of the above is a height 2 nilpotent orbit, whose HWG is presented in [3], while the second quiver is just $\mathcal{N} = 4$ QED with 2 electrons. Turning to the second and third row, we again see that for $N = 1$ the correspondence between the unitary and OSp quivers is obvious, which one may view as a further robustness of our proposal.

Now we have all the necessary ingredients to write down HWGs for the OSp quivers

$$\begin{aligned}
 \text{HWG}_{3',3'}^{(\text{I})} &= \text{PE} \left[\sum_{k=1}^N (\mu_k \mu_{2N+1-k} + \nu_k \nu_{2N+1-k}) t^{2k} \right], \\
 \text{HWG}_{3',3'}^{(\text{II})} &= \text{PE} \left[\sum_{k=1}^{N-1} (\mu_k \mu_{2N+1-k} + \nu_k \nu_{2N+1-k}) t^{2k} + (\rho^2 + \lambda^2) t^2 \right], \\
 \text{HWG}_{3',3'}^{(\text{III})} &= \text{PE} \left[\sum_{k=1}^N \mu_k \mu_{2N+1-k} t^{2k} + \sum_{j=1}^{N-1} \nu_j \nu_{2N+1-j} t^{2j} + \rho^2 t^2 \right].
 \end{aligned} \tag{3.26}$$

This proposal can be checked by a direct computation of the unrefined Hilbert series of the OSp quiver. For low values of N , results are given in table 10.

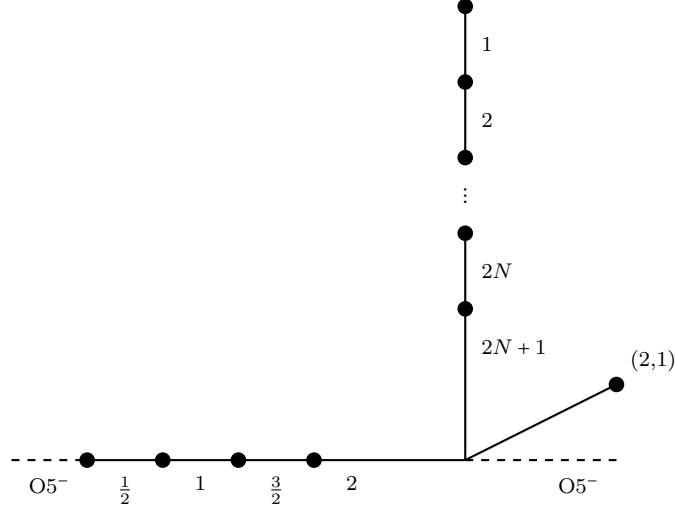


Figure 6. Orientifold web for $\text{EQ}_{3',4}$.

3.5 The $E'_3 \times E_4$ sequence

The $E'_3 \times E_4$ sequence is the magnetic quiver for the fixed point limit of the 5d IR electric quiver

$$\text{EQ}_{3',4} = \overbrace{\text{SO}(4) - \text{USp}(0) - \text{SO}(4) - \cdots - \text{USp}(0) - \text{SO}(4)}^{2N-1} \cdot \quad (3.27)$$

$$\begin{array}{c} | \\ [1\mathbf{s}] \end{array} \quad \begin{array}{c} | \\ [2\mathbf{s} + 2\mathbf{c}] \end{array}$$

It can be engineered by the orientifold web diagram presented in figure 6. Alternatively we may reformulate the electric theory $\text{EQ}_{3,4}$ as a product of two unitary electric quivers

$$\text{EQ}_{3',4} = \text{EQ}_{3'} \times \text{EQ}_4 = \overbrace{\text{SU}(2) - \text{SU}(2) - \cdots - \text{SU}(2)}^N \times \overbrace{\text{SU}(2) - \text{SU}(2) - \cdots - \text{SU}(2)}^N, \quad (3.28)$$

$$\begin{array}{c} | \\ [2\mathbf{F}] \end{array} \quad \begin{array}{c} | \\ [1\mathbf{F}] \end{array} \quad \begin{array}{c} | \\ [2\mathbf{F}] \end{array}$$

where $\text{EQ}_{3'}$ is engineered by the unitary web in figure 4, while the unitary web engineering EQ_4 is the one in figure 7. Given these webs, the magnetic quivers can be extracted using the rules in [2]. The Higgs branch of $\text{EQ}_{3',4}$ at the fixed point is the union of two cones, whose magnetic quivers are given in table 5. The HWG that we propose for the OSP quivers are

$$\text{HWG}_{3',4}^{(1)} = \text{PE} \left[\sum_{i=1}^N \mu_i \mu_{2N+1-i} t^{2i} + (\nu^2 + 1) t^2 + \nu (\mu_N q + \mu_{N+1} q^{-1}) t^{N+1} - \nu^2 \mu_N \mu_{N+1} t^{2N+2} \right]$$

$$\times \text{PE} \left[\sum_{k=1}^N \rho_k \rho_{2N+1-k} t^{2k} \right]$$

MS	OSp	Unitary
$\text{MQ}_{3',4}^{(\text{I})}$		
$\text{MQ}_{3',4}^{(\text{II})}$		

Table 5. Magnetic quivers for the two cones of the Higgs branch of $\text{EQ}_{3',4}$.

$$\begin{aligned}
 \text{HWG}_{3',4}^{(\text{II})} = & \text{PE} \left[\sum_{i=1}^N \mu_i \mu_{2N+1-i} t^{2i} + (\nu^2 + 1) t^2 + \nu (\mu_N q + \mu_{N+1} q^{-1}) t^{N+1} - \nu^2 \mu_N \mu_{N+1} t^{2N+2} \right] \\
 & \times \text{PE} \left[\sum_{k=1}^N \lambda_k \lambda_{2N+1-k} t^{2k} \right] \text{PE} [\eta^2 t^2]
 \end{aligned} \quad (3.29)$$

As a check, we have also computed the unrefined Hilbert series for low values of N which are given in table 10.

3.6 The $E_4 \times E_4$ sequence

The $E_4 \times E_4$ sequence are the magnetic quivers for the fixed point limit of the 5d electric quiver

$$\text{EQ}_{4,4} = \overbrace{[1s + 1c] - \text{SO}(4) - \text{USp}(0) - \text{SO}(4) - \dots - \text{USp}(0) - \text{SO}(4) - [2s + 2c]}^{2N-1} \quad (3.30)$$

We present the orientifold web that engineers this theory in figure 7. Alternatively we can write $\text{EQ}_{4,4}$ as the product of two unitary electric quivers

$$\text{EQ}_{4,4} = \text{EQ}_4^2 = \left(\overbrace{\text{SU}(2) - \text{SU}(2) - \dots - \text{SU}(2)}^N \right)^2, \quad (3.31)$$

\downarrow
 $[1\mathbf{F}]$

\downarrow
 $[2\mathbf{F}]$

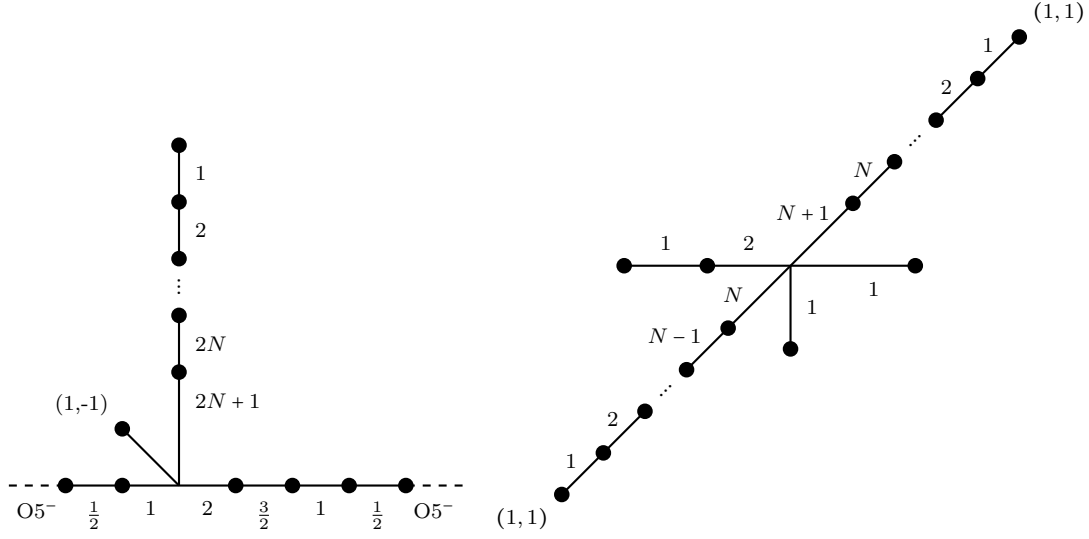


Figure 7. Brane webs for $EQ_{4,4}$ (left) and EQ_4 (right).

where each copy of EQ_4 can be engineered by the unitary web depicted in figure 7. Given the brane webs in figure 7, one can derive an OSp and a pair of unitary magnetic quivers whose Coulomb branches are expected to describe the same unique 5d Higgs branch, leading us to conjecture that

$$MQ_{4,4} = \begin{array}{c} \text{Diagram of } MQ_{4,4} \end{array} = \left(\begin{array}{c} \text{Diagram of unitary web} \end{array} \right)^2 \quad (3.32)$$

The HWG for the unitary quiver appearing here was evaluated in [34]. We can use this result to obtain an exact HWG for the OSp quiver by simply taking its square. Our claim is

$$\begin{aligned} \text{HWG}_{4,4} = & \text{PE} \left[\sum_{i=1}^N \mu_i \mu_{2N+1-i} t^{2i} + (\nu^2 + 1) t^2 + \nu (\mu_N q + \mu_{N+1} q^{-1}) t^{N+1} - \nu^2 \mu_N \mu_{N+1} t^{2N+2} \right] \\ & \times \text{PE} \left[\sum_{i=1}^N \eta_i \eta_{2N+1-i} t^{2i} + (\lambda^2 + 1) t^2 + \lambda (\eta_N r + \eta_{N+1} r^{-1}) t^{N+1} - \lambda^2 \eta_N \eta_{N+1} t^{2N+2} \right]. \end{aligned} \quad (3.33)$$

The explicit unrefined Hilbert series computation for $N = 1$ and $N = 2$ for the OSp quiver is given in table 10.

3.7 The $E_5 \times E_5$ sequence

The $E_5 \times E_5$ sequence is obtained by taking the fixed point limit of the electric quiver given by

$$EQ_{5,5} = [2s + 2c] - \overbrace{SO(4) - USp(0) - SO(4) - \dots - USp(0) - SO(4)}^{2N-1} - [2s + 2c], \quad (3.34)$$

which can be engineered using the following orientifold web:

$$(3.35)$$

Alternatively, we can rewrite $\text{EQ}_{5,5}$ as the product of two unitary electric quivers

$$\text{EQ}_{5,5} = \text{EQ}_5^2 = \left(\begin{array}{c} \overbrace{\text{SU}(2) - \text{SU}(2) - \dots - \text{SU}(2)}^N \\ | \\ [2\mathbf{F}] \end{array} \right)^2, \quad (3.36)$$

each of which is engineered by taking one copy of the following unitary brane web

$$(3.37)$$

From the brane webs in (3.35) and (3.37), we obtain the two corresponding magnetic quivers which then imply

$$\text{MQ}_{5,5} = \text{MQ}_5^2 =$$

$$= \left(\begin{array}{ccccc} & 1 & & 2 & & 1 \\ & \circ & - & \circ & - & \circ \\ & & & | & & \\ \circ & - & \dots & - & \circ & - & \dots & - & \circ \\ & 1 & & N+1 & & 1 \end{array} \right)^2. \quad (3.38)$$

The unitary quiver appearing here has been studied previously, and its HWG was given in [34]. We can now obtain the HWG for the OSp quiver by simply squaring that expression to obtain

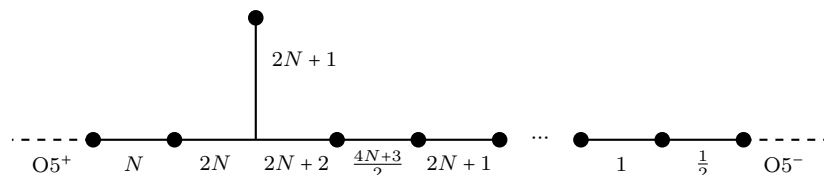
$$\begin{aligned} \text{HWG}_{5,5} = & \text{PE} \left[\sum_{i=1}^{N+1} \mu_i \mu_{2N+2-i} t^{2i} + (\nu_1^2 + \nu_2^2) t^2 + t^4 + \nu_1 \nu_2 \mu_{N+1} (t^{N+1} + t^{N+3}) - \nu_1^2 \nu_2^2 \mu_{N+1}^2 t^{2N+6} \right] \\ & \times \text{PE} \left[\sum_{i=1}^{N+1} \lambda_i \lambda_{2N+2-i} t^{2i} + (\rho_1^2 + \rho_2^2) t^2 + t^4 + \rho_1 \rho_2 \lambda_{N+1} (t^{N+1} + t^{N+3}) - \rho_1^2 \rho_2^2 \lambda_{N+1}^2 t^{2N+6} \right]. \end{aligned} \quad (3.39)$$

3.8 The $E'_5 \times E'_5$ sequence

The $E_{5'} \times E_{5'}$ sequence is obtained by considering the fixed point limit of the following 5d electric quiver:

$$\text{EQ}_{5',5'} = \overbrace{\text{SO}(4) - \text{USp}(0) - \text{SO}(4) - \dots - \text{USp}(0) - \text{SO}(4)}^{2N-1} - [4\mathbf{s} + 4\mathbf{c}], \quad (3.40)$$

which can be engineered using the orientifold web:

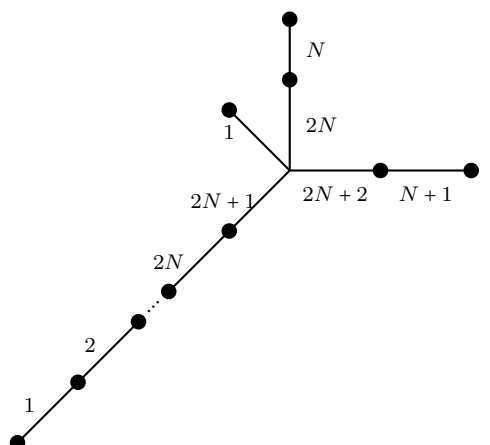


$$\quad (3.41)$$

Alternatively we may rewrite $\text{EQ}_{5',5'}$ as two copies of a single electric quiver with $\text{SU}(2)$ gauge nodes, namely

$$\text{EQ}_{5'}^2 = \left(\overbrace{\text{SU}(2) - \text{SU}(2) - \dots - \text{SU}(2)}^N \begin{array}{c} | \\ [4\mathbf{F}] \end{array} \right)^2, \quad (3.42)$$

each copy of which can now be engineered using an ordinary brane web:



$$\quad (3.43)$$

This leads us to conjecture the equivalence of the following 3d magnetic quivers

$$\begin{aligned}
 \text{MQ}_{5',5'} = & \begin{array}{c} \text{Diagram 1: A linear chain of nodes. From left to right: a red circle labeled 1, a blue circle labeled 2N, a red circle labeled 4N+1, a blue circle labeled 4N, a red circle labeled 4N+1, a blue circle labeled 4N, a red circle labeled 4N, followed by an ellipsis, a blue circle labeled 2, and a red circle labeled 2. Above the 4N+1 node is a vertical line to a red square labeled 1. Above the 4N node is a vertical line to a red square labeled 1.} \end{array} \\
 = \text{MQ}_{5'}^2 = & \left(\begin{array}{c} \text{Diagram 2: A linear chain of white circles. From left to right: a circle labeled N, a circle labeled 2N, a circle labeled 2N, followed by an ellipsis, and a circle labeled 1. Above the 2N node is a vertical line to a white circle labeled N. Above the next 2N node is a vertical line to a white circle labeled 1.} \end{array} \right)^2. \quad (3.44)
 \end{aligned}$$

We now want to write down the HWG for the Coulomb branch Hilbert series of $\text{MQ}_{5',5'}$, via the conjectured relation to the unitary quiver. The HWG for the unitary quiver $\text{MQ}_{5'}$ can be found in [35] (also see table 18 of [36]), which we report in table 9. Consequently, the HWG for $\text{MQ}_{5',5'}$ is obtained by squaring this result, namely

$$\text{HWG}_{5',5'} = \text{PE} \left[\sum_{k=1}^N (\mu_{2k} + \nu_{2k}) t^{2k} \right]. \quad (3.45)$$

3.9 The $E_{5'} \times E_6$ sequence

The $E_{5'} \times E_6$ sequence is obtained by taking the fixed point limit of the 5d electric quiver given by

$$\text{EQ}_{5',6} = \overbrace{[1\mathbf{s}] - \text{SO}(4) - \text{USp}(0) - \text{SO}(4) - \cdots - \text{USp}(0) - \text{SO}(4) - [4\mathbf{s} + 4\mathbf{c}]}^{2N-1}, \quad (3.46)$$

which can be engineered by the following orientifold web diagram

$$\begin{array}{c} \text{Diagram 3: An orientifold web diagram. A vertical line of black dots is on the left. From top to bottom: a dot labeled 1, a dot labeled 2, an ellipsis, a dot labeled 2N, and a dot labeled 2N+1. A diagonal line connects the 2N+1 dot to a dot on a horizontal line. This horizontal line has dots labeled O5-, 2N+1/2, 2N+1, 2N+2, 4N+3/2, 1, 1/2, and O5-. A label (1,-1) is near the 2N+1 dot.} \end{array} \quad (3.47)$$

Alternatively we may reformulate $\text{EQ}_{5',6}$ as the product of two unitary electric quivers

$$\text{EQ}_{5',6} = \text{EQ}_{5'} \times \text{EQ}_6 = \overbrace{\text{SU}(2) - \text{SU}(2) - \cdots - \text{SU}(2)}^N \times \overbrace{\text{SU}(2) - \text{SU}(2) - \cdots - \text{SU}(2)}^N, \quad (3.48)$$

\downarrow
 $[4\mathbf{F}]$

\downarrow
 $[1\mathbf{F}]$

\downarrow
 $[4\mathbf{F}]$

where $\text{EQ}_{5'}$ is engineered by the unitary web in (3.43), and EQ_6 is the IR quiver description of the web diagram given by

(3.49)

Reading off the OSp magnetic quiver from (3.47) and the unitary quivers from (3.43), (3.49), we arrive at the conjecture

$$\begin{aligned} \text{MQ}_{5',6} &= \text{Diagram 1} = \text{Diagram 2} \\ &= \text{MQ}_{5'} \times \text{MQ}_6 = \text{Diagram 3} \times \text{Diagram 4} \end{aligned} \quad (3.50)$$

Diagram 1: A chain of nodes with labels 2N, 4N+2, 4N+2, 4N+2, ..., 2, 2. The second node is blue, the third is red, and the last is red.

Diagram 2: A chain of nodes with labels N, 2N, 2N, ..., 1. The second node has a vertical line to a blue node labeled 2N, and the third has a vertical line to a white node labeled 1.

Diagram 3: A chain of nodes with labels 1, N+1, 2N+1, ..., 1. The third node has a vertical line to a white node labeled 1.

Diagram 4: A chain of nodes with labels 1, N+1, 2N+1, ..., 1. The third node has a vertical line to a white node labeled N+1.

Obtaining the HWG for the Coulomb branch of $\text{MQ}_{5',6}$ is now straightforward, given the above relation. The HWG for MQ'_5 was worked out in [35, 36] and the HWG for MQ_6 appears in [34]. For the OSp quiver $\text{MQ}_{5',6}$ in (3.50) we simply take the product of these two results

$$\text{HWG}_{5',6} = \text{PE} \left[\sum_{i=1}^N (\mu_{2i} + \nu_{2i}) t^{2i} + t^2 + (\mu_{2N+2} q + \mu_{2N+3} q^{-1}) t^{N+1} \right]. \quad (3.51)$$

This can be verified by computing the unrefined Coulomb branch Hilbert series, the results of which are given in table 10 for low values of N .

3.10 The $E_6 \times E_6$ sequence

The electric quiver for the IR limit of the $E_6 \times E_6$ sequence reads

$$\text{EQ}_{6,6} = [1\mathbf{s} + 1\mathbf{c}] - \overbrace{\text{SO}(4) - \text{USp}(0) - \text{SO}(4) - \cdots - \text{USp}(0) - \text{SO}(4)}^{2N-1} - [4\mathbf{s} + 4\mathbf{c}] \quad (3.52)$$

It can be engineered using the following orientifold web

$$\quad (3.53)$$

Alternatively we may present the electric quiver as a product of two unitary quivers:

$$\text{EQ}_{6,6} = \text{EQ}_6^2 = \left(\overbrace{\text{SU}(2) - \text{SU}(2) - \cdots - \text{SU}(2)}^N \begin{array}{c} [1\mathbf{F}] \\ [4\mathbf{F}] \end{array} \right)^2, \quad (3.54)$$

where each copy is engineered by the web diagram in (3.49). We can then obtain the magnetic quivers from the brane webs in (3.53) and (3.49), that lead us to the conjecture that

$$\text{MQ}_{6,6} = \text{MQ}_6^2 = \left(\begin{array}{c} \text{web diagram} \\ \text{1} \\ \text{N+1} \\ \text{2N+1} \end{array} \right)^2. \quad (3.55)$$

This has an immediate corollary which allows us to extract the HWG for the OSp quiver appearing above by squaring the known result [34] for the unitary quiver:

$$\begin{aligned} \text{HWG}_{6,6} = \text{PE} & \left[\sum_{i=1}^N \mu_{2i} t^{2i} + t^2 + (\mu_{2N+2} q_1 + \mu_{2N+3} q_1^{-1}) t^{N+1} \right] \\ & \times \text{PE} \left[\sum_{i=1}^N \nu_{2i} t^{2i} + t^2 + (\nu_{2N+2} q_2 + \nu_{2N+3} q_2^{-1}) t^{N+1} \right]. \end{aligned} \quad (3.56)$$

Nahm partition	$(\delta n_h, \delta n_v)$
$[1^8]$	$(112, 100)$
$[3^2, 1^2]$	$(72, 69)$
$[3^2, 1^2]$	$(72, 69)$

Table 6. Table containing the data defining the punctures for the 4d $E_6 \times E_6$ theory.

Indeed the unrefined Hilbert series for the OSp quiver for $N = 1$ was computed in [6], and is in agreement with our claim. For higher values of N we were not able to perform an explicit computation due to the high rank of the OSp quiver. This is one instance in which our conjecture proves powerful, as it gives an exact expression for the Hilbert series of a quiver which would otherwise be very challenging to compute.

We further point out an interesting fact about this theory, namely the existence of a 4d $\mathcal{N} = 2$ theory with very similar properties. As noticed in [26], there is one class-S theory of D_4 type in which a single three-punctured sphere realizes a product SCFT, where both factors are the E_6 Minahan-Nemeschansky (MN) theory [37]. We recall that the E_6 Minahan-Nemeschansky theory is a 4d $\mathcal{N} = 2$ SCFT of rank 1, with flavor symmetry group E_6 , and central charges

$$a_{E_6} = \frac{41}{24}, \quad c_{E_6} = \frac{13}{6}, \quad (3.57)$$

We report the partitions labeling the punctures in table 6, together with their contribution to the effective number of hypermultiplets and vector multiplets.

From this data it is easy to compute the central charges a and c of this theory,⁶ finding

$$a_{E_6 \times E_6} = \frac{41}{12}, \quad c_{E_6 \times E_6} = \frac{13}{3} \quad (3.58)$$

as it should be for two copies of the E_6 MN theory. By applying the procedure to write the 3d mirror for this theory⁷ we find that the full puncture $[1^8]$ is associated to the quiver tail (3.59) while the puncture $[3^2, 1^2]$ is associated to the quiver tail (3.60),

$$\begin{array}{ccccccc} \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \blacksquare \\ 2 & 2 & 4 & 4 & 6 & 6 & 8 \end{array} \quad (3.59)$$

$$\begin{array}{ccc} \bullet & \bullet & \blacksquare \\ 2 & 4 & 8 \end{array} . \quad (3.60)$$

Gluing the three tails together results in the magnetic quiver for the 5d $E_6 \times E_6$ depicted in (3.55). Therefore the magnetic quiver of the 5d $E_6 \times E_6$ theory is the 3d mirror theory of the 4d $E_6 \times E_6$ theory above described. It is then tempting to conjecture that the 5d $E_6 \times E_6$ theory reduces to 4d to this D_4 type class-S theory, giving two copies of E_6 Minahan-Nemeschansky.

Having derived the magnetic quiver for the $E_6 \times E_6$ sequence from the brane web, for any $N \in \mathbb{N}$, we can use the same argument as the paragraphs above to conjecture that all

⁶For a small review see appendix D.1.

⁷For a small review see appendix D.2.

class-S theories of D_{2N+2} type given by a three punctured sphere with regular punctures given by $[1^{4N+4}], [2N+1, 2N+1, 1, 1], [2N+1, 2N+1, 1, 1]$ will be a factorized SCFT. We conjecture that it will decompose into two copies of three punctured A_{2N} spheres, with regular punctures given by $[1^{2N+1}], [N^2, 1], [N^2, 1]$. It will be interesting to further check this proposal.

3.11 The $E_7 \times E_7$ sequence

The $E_7 \times E_7$ sequence corresponds to the fixed point limit of the following IR quiver

$$\text{EQ}_{7,7} = [2\mathbf{s} + 2\mathbf{c}] - \overbrace{\text{SO}(4) - \text{USp}(0) - \text{SO}(4) - \cdots - \text{USp}(0) - \text{SO}(4)}^{2N-1} - [4\mathbf{s} + 4\mathbf{c}] \quad . \quad (3.61)$$

It can be engineered using the orientifold web given by

$$\quad \quad \quad (3.62)$$

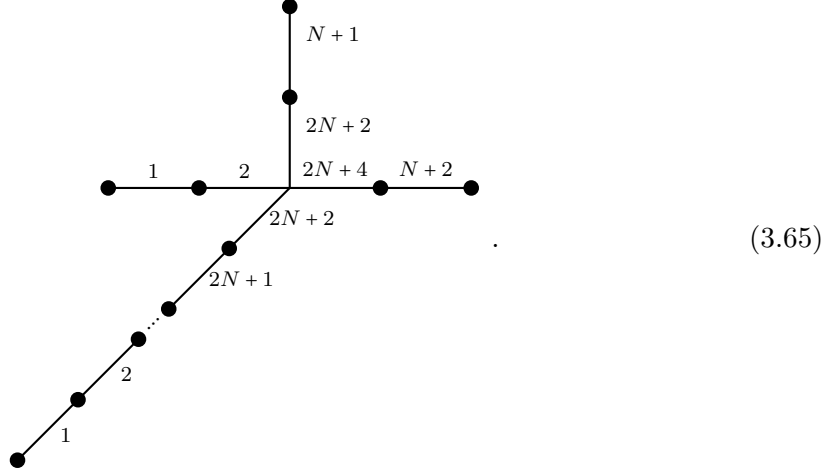
This web can be converted to a magnetic quiver following [2], which results in

$$\text{MQ}_{7,7} = \quad \quad \quad (3.63)$$

Now we use the alternative description of the $\text{EQ}_{7,7}$ as a product of a pair of linear quivers with $\text{SU}(2)$ nodes, namely

$$\text{EQ}_{7,7} = \text{EQ}_7^2 = \left(\overbrace{\text{SU}(2) - \text{SU}(2) - \cdots - \text{SU}(2)}^N \begin{array}{c} | \\ [2\mathbf{F}] \end{array} \quad \begin{array}{c} | \\ [4\mathbf{F}] \end{array} \right)^2 \quad . \quad (3.64)$$

Each individual factor can be engineered using the following unitary 5-brane web



The magnetic quiver one obtains from this unitary web leads us to the following conjecture

$$\text{MQ}_{7,7} = \text{MQ}_7^2 = \left(\begin{array}{ccccccc} & & & N+1 & & & \\ & & & \circ & & & \\ & & & | & & & \\ \circ & \circ & \circ & \circ & \circ & \dots & \circ \\ 1 & 2 & N+2 & 2N+2 & 2N+1 & & 1 \end{array} \right)^2. \quad (3.66)$$

The HWG for the unitary quiver appearing above was conjectured in [34]. We will use this result and square it to obtain the HWG for the OSp quiver $\text{MQ}_{7,7}$ (3.63):

$$\begin{aligned} \text{HWG}_{7,7} = & \text{PE} \left[\sum_{i=1}^{N+1} \mu_{2i} t^{2i} + t^4 + \nu^2 t^2 + \nu \mu_{2N+4} (t^{N+1} + t^{N+3}) + \mu_{2N+4}^2 t^{2N+4} - \nu^2 \mu_{2N+4} t^{2N+6} \right] \\ & \times \text{PE} \left[\sum_{i=1}^{N+1} \lambda_{2i} t^{2i} + t^4 + \rho^2 t^2 + \rho \lambda_{2N+4} (t^{N+1} + t^{N+3}) + \lambda_{2N+4}^2 t^{2N+4} - \rho^2 \lambda_{2N+4} t^{2N+6} \right]. \end{aligned} \quad (3.67)$$

We further point out an interesting fact about this theory, namely the existence of a 4d $\mathcal{N} = 2$ theory with very similar properties. As noticed in [26], there is one class-S theory of D_5 type in which a single three-punctured sphere realizes a product SCFT, where both factors are the E_7 Minahan-Nemeschansky theory [38]. We recall that the E_7 Minahan-Nemeschansky theory is a 4d $\mathcal{N} = 2$ SCFT of rank 1, with flavor symmetry group E_7 , and central charges

$$a_{E_7} = \frac{59}{24}, \quad c_{E_7} = \frac{19}{6}. \quad (3.68)$$

We report the partitions labeling the punctures in table 7, together with their contribution to the effective number of hypermultiplets and vector multiplets.

From this data it is easy to compute the central charges a and c of this theory, finding

$$a_{E_7 \times E_7} = \frac{59}{12}, \quad c_{E_7 \times E_7} = \frac{19}{3} \quad (3.69)$$

as it should be for two copies of the E_7 MN theory.

Nahm partition	$(\delta n_h, \delta n_v)$
$[1^{10}]$	$(240, 220)$
$[5^2]$	$(104, 102)$
$[3^2, 1^4]$	$(184, 177)$

Table 7. Table containing the data defining the punctures for the 4d $E_7 \times E_7$ theory.

By applying the procedure to write the 3d mirror for this theory we find that the full puncture $[1^{10}]$ is associated to the quiver tail (3.70), the puncture $[3^2, 1^4]$ is associated to the quiver tail (3.71), and the puncture $[5^2]$ is associated to the quiver tail (3.72),

$$\begin{array}{ccccccccccc} \bullet & & \bullet & & \bullet & & \bullet & & \bullet & & \blacksquare \\ 2 & & 2 & & 4 & & 4 & & 6 & & 6 & & 8 & & 8 & & 10 \end{array} \quad (3.70)$$

$$\begin{array}{ccccccc} \bullet & & \bullet & & \bullet & & \bullet & & \blacksquare \\ 2 & & 2 & & 4 & & 8 & & 10 \end{array} \quad (3.71)$$

$$\begin{array}{ccc} \bullet & & \blacksquare \\ 4 & & 10 \end{array} . \quad (3.72)$$

Gluing the three tails together results in the magnetic quiver for the 5d $E_7 \times E_7$ depicted in (3.63) for $N = 1$. Therefore the magnetic quiver of the 5d $E_7 \times E_7$ theory is the 3d mirror theory of the 4d $E_7 \times E_7$ theory above described. It is then tempting to conjecture that the 5d $E_7 \times E_7$ theory reduces to 4d to this D_5 type class-S theory, giving two copies of E_7 Minahan-Nemeschansky.

Having derived the magnetic quiver for the $E_7 \times E_7$ sequence from the brane web, for any $N \in \mathbb{N}$, we can use the same argument as the paragraphs above to conjecture that all class-S theories of D_{2N+3} type given by a three punctured sphere with regular punctures given by $[1^{4N+6}], [2N+1, 2N+1, 1^4], [2N+3, 2N+3]$ will be a factorized SCFT. We conjecture that it will decompose into two copies of three punctured A_{2N+1} spheres, with regular punctures given by $[1^{2N+2}], [N+1, N+1], [N, N, 1, 1]$. It will be interesting to further check this proposal.

3.12 An outlier: the $E_8 \times E_8$ theory

While not explicitly written⁸ in [26], it is easy to use the methods of such paper to find a choice of punctures in the D_6 theory, such that we realize the product of two copies of the E_8 Minahan-Nemeschansky theory [38]. We recall that the E_8 Minahan-Nemeschansky theory is a 4d $\mathcal{N} = 2$ SCFT of rank 1, with flavor symmetry group E_8 , and central charges

$$a_{E_8} = \frac{95}{24}, \quad c_{E_8} = \frac{31}{6}. \quad (3.73)$$

We report the partitions labeling the punctures which we believe engineer this product SCFT in table 8, together with their contribution to the effective number of hypermultiplets and vector multiplets.

⁸But surely noticed by the authors of such paper. See for example [27] and [28] for discussions about product SCFTs in class-S.

4 Discussion

In this paper we studied 3d $\mathcal{N} = 4$ OSp quivers whose moduli space of vacua is comprised of two decoupled sectors. These OSp quivers were derived, using brane webs with O5-planes, as the magnetic quivers for the infinite gauge coupling limit of 5d $\mathcal{N} = 1$ gauge theories whose gauge group is a product of $SO(4)$ factors and contain matter hypermultiplets transforming as spinors of opposite chirality under the gauge group factors. We argued for the proposed factorisation, exploiting an accidental isomorphism between the Lie algebra of $SO(4)$ and $SU(2) \times SU(2)$, by rewriting the gauge theory in terms of the group $SU(2)$ and then taking the infinite coupling limit on both sides. The resulting theory in terms of $SU(2)$ gauge groups is generically comprised of two decoupled sectors, each of which we also engineered using ordinary brane webs, without O5-planes. The ordinary brane webs were subsequently used to derive unitary magnetic quivers, which we then used to propose as the components to which the OSp quivers factorise. We further used this correspondence to extract highest weight generating functions for the Coulomb branch Hilbert series of the OSp quivers, relying on existing results for the unitary quivers. In some cases where the unitary quivers had not previously appeared in the literature, we also computed the highest weight generators. In order to test our proposal for the factorisation, and consequently the conjectured highest weight generators for the OSp quivers, we also computed the unrefined Coulomb branch Hilbert series of the OSp quivers directly in a perturbative manner and found an agreement with the proposed HWGs. We further illustrated the matching of the Higgs branch Hilbert series in two cases where we were able to perform the computation exactly on the OSp side. These too were in agreement with the results of the Higgs branch Hilbert series of the unitary side.

Although the higher dimensional intuition has led us to derive these results. It begs the question, whether a truly three-dimensional logic can be used to argue for or provide an explanation for the factorisation property of these OSp quivers. Moreover, all the quantitative checks performed in this paper probe the moduli space of the quivers studied. It would be interesting to ask whether the relationship between the OSp and unitary quivers in this paper are full-fledged dualities, or just a formal relation between their moduli spaces of vacua. One potential check that can be performed to illuminate this question would be to compute other observables, such as the superconformal index, or the three-sphere partition function for the theories in question.

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A Higgs branch of $E_1 \times E_1$ sequence

In this appendix we give a derivation of the formula (3.7), for the Higgs branch Hilbert series of $\text{MQ}_{1,1}$ (3.3). The Higgs branch of the quiver (3.3) can be computed using the gluing technique, following the discussion in appendix A of [39]. The first step is to break down the OSp quiver in (3.3) into pieces as follows

$$\begin{array}{ccccccc} \bigcirc & \cdots & \bigcirc & \square & \square & \blacksquare & \blacksquare \text{---} \blacksquare \\ 1 & & 2N-2 & 2N-1 & 2N-1 & 2N & 2N \quad 4 \end{array} \quad (\text{A.1})$$

The Higgs Branch Hilbert series of the original quiver (3.3) is then obtained by taking the product of the Hilbert series of the individual factors above, together with the gluing factors associated with the $\text{U}(2N-1)$ and $\text{USp}(2N)$ nodes which are to be gauged, all integrated with the appropriate Haar measure for the aforementioned gauge groups

$$\int d\mu_{\text{U}(2N-1)} \int d\mu_{C_N} H_{T[\text{SU}(2N-1)]} H_{\text{glue}}^{(2N-1)} H_{[2N-1]-[C_N]} H_{\text{glue}}^{(C_N)} H_{[C_N]-[D_2]}. \quad (\text{A.2})$$

The individual pieces in the above integrand are as follows

$$\begin{aligned} H_{T[\text{SU}(2N-1)]} &= \prod_{q=2}^{2N-1} (1-t^{2q}) \text{PE} \left[\chi_{[1,0,\dots,0,1]}^{\text{SU}(2N-1)} t^2 \right], & H_{[C_N]-[D_2]} &= \text{PE} \left[4\chi_{[1,0,\dots,0]}^{C_N} t \right], \\ H_{[2N-1]-[C_N]} &= \text{PE} \left[\left(\chi_{[1,0,\dots,0]}^{\text{SU}(2N-1)} q + \chi_{[0,\dots,0,1]}^{\text{SU}(2N-1)} q^{-1} \right) \chi_{[1,0,\dots,0]}^{C_N} t \right], \\ H_{\text{glue}}^{(2N-1)} &= \frac{(1-t^2)}{\text{PE} \left[\chi_{[1,0,\dots,0,1]}^{\text{SU}(2N-1)} t^2 \right]}, & H_{\text{glue}}^{(C_N)} &= \frac{1}{\text{PE} \left[\chi_{[2,0,\dots,0]}^{C_N} t^2 \right]}. \end{aligned} \quad (\text{A.3})$$

Plugging these into the integral (A.2) we arrive at

$$\prod_{q=1}^{2N-1} (1-t^{2q}) \int d\mu_{\text{U}(2N-1)} \int d\mu_{C_N} H_{[2N-1]-[C_N]} H_{\text{glue}}^{(C_N)} H_{[C_N]-[D_2]}. \quad (\text{A.4})$$

Next we perform the integral over the $\text{U}(2N-1)$ group, by counting the gauge invariants of the free theory $[2N-1] - [C_N]$, which is the only part of the integrand that sees this integral (see appendix A of [39] for more details)

$$\int d\mu_{\text{U}(2N-1)} H_{[2N-1]-[C_N]} = (1-t^{4N}) \text{PE} \left[\left(1 + \chi_{[2,0,\dots,0]}^{C_N} + \chi_{[0,1,\dots,0]}^{C_N} \right) t^2 \right]. \quad (\text{A.5})$$

Substituting this back into (A.4) one obtains the desired formula (3.7) for the Higgs branch Hilbert series of $\text{MQ}_{1,1}$ (3.3).

B List of HWG of unitary quivers

In this appendix we tabulate the highest weight generating functions of all the unitary quivers appearing in the previous sections.

Table 9. List of the HWG for the unitary quivers appearing in the earlier sections. The corresponding fugacities are denoted by subscripts in the symmetry groups. Note that for $N = 1$, there is a possible enhancement in the symmetry which can be read from the set of balanced nodes in the quivers for $N = 1$.

Quiver	Symmetry	PL[HWG]
	$SU(2N)_\mu$	$\sum_{k=1}^N \mu_k \mu_{2N-k} t^{2k}$
	$SU(2N)_\mu \times U(1)_q$	$t^2 + (q + q^{-1}) \mu_N t^{N+1} + \sum_{k=1}^N \mu_k \mu_{2N-k} t^{2k} - \mu_N^2 t^{2N+2}$
	$SU(2N)_\mu \times U(1)_q$	$t^2 + (\mu_{N+1} q + \mu_{N-1} q^{-1}) t^{N+1} + \sum_{k=1}^{N-1} \mu_k \mu_{2N-k} t^{2k} - \mu_{N+1} \mu_{N-1} t^{2N+2}$
	$SU(2N+1)_\mu$	$\sum_{k=1}^N \mu_k \mu_{2N+1-k} t^{2k}$

	$\text{SU}(2N+1)_\mu \times \text{SU}(2)_\nu \times \text{U}(1)_q$	$\sum_{k=1}^N \mu_k \mu_{2N+1-k} t^{2k} + (\nu^2 + 1)t^2 + \nu(\mu_N q + \mu_{N+1} q^{-1})t^{N+1} - \nu^2 \mu_N \mu_{N+1} t^{2N+2}$
	$\text{SU}(2N+1)_\mu \times \text{SU}(2)_\nu$	$\sum_{k=1}^{N-1} \mu_k \mu_{2N+1-k} t^{2k} + \nu^2 t^2$
	$\text{SU}(2N+2)_\mu \times \text{SU}(2)_{\nu_1} \times \text{SU}(2)_{\nu_2}$	$\sum_{k=1}^{N+1} \mu_k \mu_{2N+2-k} t^{2k} + (\nu_1^2 + \nu_2^2)t^2 + t^4 + \nu_1 \nu_2 \mu_{N+1} (t^{N+1} + t^{N+3}) - \nu_1^2 \nu_2^2 \mu_{N+1}^2 t^{2N+6}$
	$\text{SO}(4N+6)_\mu$	$\sum_{k=1}^N \mu_{2k} t^{2k}$
	$\text{SO}(4N+6)_\mu \times \text{U}(1)_q$	$\sum_{k=1}^N \mu_{2k} t^{2k} + (\mu_{2N+2} q + \mu_{2N+3} q^{-1}) t^{N+1} + t^2$
	$\text{SO}(4N+8)_\mu \times \text{SU}(2)_\nu$	$\sum_{k=1}^{N+1} \mu_{2k} t^{2k} + t^4 + \nu \mu_{2N+4} (t^{N+1} + t^{N+3}) + \mu_{2N+4}^2 t^{2N+4} + \nu^2 t^2 - \nu^2 \mu_{2N+4} t^{2N+6}$

C Unrefined Coulomb branch Hilbert series for low rank OSp quivers

In this appendix, we quote the results for the unrefined Hilbert series computed for the orthosymplectic quivers. These results match with the perturbative results obtained from the HWG listed in the table 1.

Table 10. The computation of Coulomb branch Hilbert series of orthosymplectic quivers given in main sections for small values of N . The full Hilbert series is given by summing over the integer and half-integer sublattices of the magnetic weights.

Quiver	Coulomb branch Hilbert series		
	$\vec{m} \in \mathbb{Z}$	$\vec{m} \in \mathbb{Z} + 1/2$	HS
$\text{MQ}_{1,1} _{N=2}$	$1 + 30t^2 + 433t^4 + 4070t^6 + 28384t^8 + 158174t^{10} + \mathcal{O}(t^{12})$	not required	$1 + 30t^2 + 433t^4 + 4070t^6 + 28384t^8 + 158174t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{1,1} _{N=3}$	$1 + 70t^2 + 2413t^4 + 54670t^6 + 917244t^8 + 12178110t^{10} + \mathcal{O}(t^{12})$	not required	$1 + 70t^2 + 2413t^4 + 54670t^6 + 917244t^8 + 12178110t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{1,3}^{(\text{I})} _{N=2}$	$1 + 31t^2 + 464t^4 + 4574t^6 + 33908t^8 + 203160t^{10} + \mathcal{O}(t^{12})$	$12t^3 + 320t^5 + 4188t^7 + 36488t^9 + \mathcal{O}(t^{11})$	$1 + 31t^2 + 12t^3 + 464t^4 + 320t^5 + 4574t^6 + 4188t^7 + 33908t^8 + 36488t^9 + 203160t^{10} + \mathcal{O}(t^{11})$
$\text{MQ}_{1,3}^{(\text{I})} _{N=3}$	$1 + 71t^2 + 2484t^4 + 57154t^6 + 974748t^8 + 13173046t^{10} + \mathcal{O}(t^{12})$	$40t^4 + 2520t^6 + 78584t^8 + 1619760t^{10} + \mathcal{O}(t^{12})$	$1 + 71t^2 + 2524t^4 + 59674t^6 + 1053332t^8 + 14792806t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{1,3}^{(\text{II})} _{N=2}$	$1 + 31t^2 + 444t^4 + 4059t^6 + 27344t^8 + 147137t^{10} + \mathcal{O}(t^{12})$	$8t^3 + 200t^5 + 2432t^7 + 19560t^9 + \mathcal{O}(t^{11})$	$1 + 31t^2 + 8t^3 + 444t^4 + 200t^5 + 4059t^6 + 2432t^7 + 27344t^8 + 19560t^9 + 147137t^{10} + \mathcal{O}(t^{11})$
$\text{MQ}_{1,3}^{(\text{II})} _{N=3}$	$1 + 71t^2 + 2484t^4 + 56979t^6 + 964339t^8 + 12865508t^{10} + \mathcal{O}(t^{12})$	$30t^4 + 1848t^6 + 56250t^8 + 1129770t^{10} + \mathcal{O}(t^{12})$	$1 + 71t^2 + 2514t^4 + 58827t^6 + 1020589t^8 + 13995278t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3,3}^{(\text{I})} _{N=1}$	$1 + 8t^2 + 62t^4 + 280t^6 + 1011t^8 + 2944t^{10} + \mathcal{O}(t^{12})$	$8t^2 + 56t^4 + 280t^6 + 992t^8 + 2944t^{10} + \mathcal{O}(t^{12})$	$1 + 16t^2 + 118t^4 + 560t^6 + 2003t^8 + 5888t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3,3}^{(\text{I})} _{N=2}$	$1 + 32t^2 + 496t^4 + 5254t^6 + 43368t^8 + 294996t^{10} + \mathcal{O}(t^{12})$	$24t^3 + 664t^5 + 9040t^7 + 82976t^9 + \mathcal{O}(t^{11})$	$1 + 32t^2 + 24t^3 + 496t^4 + 664t^5 + 5254t^6 + 9040t^7 + 43368t^8 + 82976t^9 + 294996t^{10} + \mathcal{O}(t^{11})$
$\text{MQ}_{3,3}^{(\text{II})} _{N=1}$	$1 + 2t^2 + 11t^4 + 20t^6 + 45t^8 + 70t^{10} + \mathcal{O}(t^{12})$	$4t^2 + 8t^4 + 24t^6 + 40t^8 + 76t^{10} + \mathcal{O}(t^{12})$	$1 + 6t^2 + 19t^4 + 44t^6 + 85t^8 + 146t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3,3}^{(\text{II})} _{N=2}$	$1 + 32t^2 + 456t^4 + 4104t^6 + 27490t^8 + 148792t^{10} + \mathcal{O}(t^{12})$	$16t^3 + 416t^5 + 4960t^7 + 38400t^9 + \mathcal{O}(t^{11})$	$1 + 32t^2 + 16t^3 + 456t^4 + 416t^5 + 4104t^6 + 4960t^7 + 27490t^8 + 38400t^9 + 148792t^{10} + \mathcal{O}(t^{11})$
$\text{MQ}_{3,3}^{(\text{III})} _{N=1}$	$1 + 5t^2 + 30t^4 + 94t^6 + 263t^8 + 587t^{10} + \mathcal{O}(t^{12})$	$6t^2 + 26t^4 + 98t^6 + 254t^8 + 596t^{10} + \mathcal{O}(t^{12})$	$1 + 11t^2 + 56t^4 + 192t^6 + 517t^8 + 1183t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3,3}^{(\text{III})} _{N=2}$	$1 + 32t^2 + 476t^4 + 4671t^6 + 34989t^8 + 214034t^{10} + \mathcal{O}(t^{12})$	$20t^3 + 540t^5 + 6920t^7 + 58628t^9 + \mathcal{O}(t^{11})$	$1 + 32t^2 + 20t^3 + 476t^4 + 540t^5 + 4671t^6 + 6920t^7 + 34989t^8 + 58628t^9 + 214034t^{10} + \mathcal{O}(t^{11})$
$\text{MQ}_{3',3'}^{(\text{I})} _{N=1}$	$1 + 16t^2 + 118t^4 + 560t^6 + 2003t^8 + 5888t^{10} + \mathcal{O}(t^{12})$	not required	$1 + 16t^2 + 118t^4 + 560t^6 + 2003t^8 + 5888t^{10} + \mathcal{O}(t^{12})$

$\text{MQ}_{3',3'}^{(\text{I})} _{N=2}$	$1 + 48t^2 + 1126t^4 + 17248t^6 + 194729t^8 + 1735152t^{10} + \mathcal{O}(t^{12})$	not required	$1 + 48t^2 + 1126t^4 + 17248t^6 + 194729t^8 + 1735152t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3',3'}^{(\text{II})} _{N=1}$	$1 + 6t^2 + 19t^4 + 44t^6 + 85t^8 + 146t^{10} + \mathcal{O}(t^{12})$	not required	$1 + 6t^2 + 19t^4 + 44t^6 + 85t^8 + 146t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3',3'}^{(\text{II})} _{N=2}$	$1 + 54t^2 + 1283t^4 + 18412t^6 + 185691t^8 + 1438022t^{10} + \mathcal{O}(t^{12})$	not required	$1 + 54t^2 + 1283t^4 + 18412t^6 + 185691t^8 + 1438022t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3',3'}^{(\text{III})} _{N=1}$	$1 + 11t^2 + 56t^4 + 192t^6 + 517t^8 + 1183t^{10} + \mathcal{O}(t^{12})$	not required	$1 + 11t^2 + 56t^4 + 192t^6 + 517t^8 + 1183t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3',3'}^{(\text{III})} _{N=2}$	$1 + 51t^2 + 1200t^4 + 17824t^6 + 191099t^8 + 1596553t^{10} + \mathcal{O}(t^{12})$	not required	$1 + 51t^2 + 1200t^4 + 17824t^6 + 191099t^8 + 1596553t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3',4}^{(\text{I})} _{N=1}$	$1 + 20t^2 + 227t^4 + 1720t^6 + 9552t^8 + 42168t^{10} + \mathcal{O}(t^{12})$	$12t^2 + 192t^4 + 1592t^6 + 9184t^8 + 41224t^{10} + \mathcal{O}(t^{12})$	$1 + 32t^2 + 419t^4 + 3312t^6 + 18736t^8 + 83392t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3',4}^{(\text{I})} _{N=2}$	$1 + 52t^2 + 1327t^4 + 22500t^6 + 286968t^8 + 2939292t^{10} + \mathcal{O}(t^{12})$	$40t^3 + 1780t^5 + 39140t^7 + 570120t^9 + \mathcal{O}(t^{11})$	$1 + 52t^2 + 40t^3 + 1327t^4 + 1780t^5 + 22500t^6 + 39140t^7 + 286968t^8 + 570120t^9 + 2939292t^{10} + \mathcal{O}(t^{11})$
$\text{MQ}_{3',4}^{(\text{II})} _{N=1}$	$1 + 15t^2 + 145t^4 + 879t^6 + 3964t^8 + 14388t^{10} + \mathcal{O}(t^{12})$	$12t^2 + 132t^4 + 848t^6 + 3888t^8 + 14240t^{10} + \mathcal{O}(t^{12})$	$1 + 27t^2 + 277t^4 + 1727t^6 + 7852t^8 + 28628t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{3',4}^{(\text{II})} _{N=2}$	$1 + 55t^2 + 1413t^4 + 23399t^6 + 287256t^8 + 2808576t^{10} + \mathcal{O}(t^{12})$	$40t^3 + 1900t^5 + 41680t^7 + 585260t^9 + \mathcal{O}(t^{11})$	$1 + 55t^2 + 40t^3 + 1413t^4 + 1900t^5 + 23399t^6 + 41680t^7 + 287256t^8 + 585260t^9 + 2808576t^{10} + \mathcal{O}(t^{11})$
$\text{MQ}_{4,4} _{N=1}$	$1 + 24t^2 + 496t^4 + 5800t^6 + 47734t^8 + 299176t^{10} + \mathcal{O}(t^{12})$	$24t^2 + 480t^4 + 5800t^6 + 47616t^8 + 299176t^{10} + \mathcal{O}(t^{12})$	$1 + 48t^2 + 976t^4 + 11600t^6 + 95350t^8 + 598352t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{4,4} _{N=2}$	$1 + 56t^2 + 1544t^4 + 30192t^6 + 468888t^8 + 5964152t^{10} + \mathcal{O}(t^{12})$	$80t^3 + 3880t^5 + 93240t^7 + 1510680t^9 + \mathcal{O}(t^{11})$	$1 + 56t^2 + 80t^3 + 1544t^4 + 3880t^5 + 30192t^6 + 93240t^7 + 468888t^8 + 1510680t^9 + 5964152t^{10} + \mathcal{O}(t^{11})$
$\text{MQ}_{5,5} _{N=1}$	$1 + 42t^2 + 1805t^4 + 42204t^6 + 693740t^8 + 8548816t^{10} + \mathcal{O}(t^{12})$	$48t^2 + 1760t^4 + 42384t^6 + 692960t^8 + 8551232t^{10} + \mathcal{O}(t^{12})$	$1 + 90t^2 + 3565t^4 + 84588t^6 + 1386700t^8 + 17100048t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{5,5} _{N=2}$	$1 + 82t^2 + 3329t^4 + \mathcal{O}(t^6)$	$160t^3 + \mathcal{O}(t^5)$	$1 + 82t^2 + 160t^3 + 3329t^4 + \mathcal{O}(t^5)$
$\text{MQ}_{5',5'} _{N=1}$	$1 + 90t^2 + 3565t^4 + 84588t^6 + 1386700t^8 + 17100048t^{10} + \mathcal{O}(t^{12})$	not required	$1 + 90t^2 + 3565t^4 + 84588t^6 + 1386700t^8 + 17100048t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{5',5'} _{N=2}$	$1 + 182t^2 + 16443t^4 + 977366t^6 + \mathcal{O}(t^8)$	not required	$1 + 182t^2 + 16443t^4 + 977366t^6 + \mathcal{O}(t^8)$
$\text{MQ}_{5',6} _{N=1}$	$1 + 91t^2 + 4118t^4 + 122828t^6 + 2660832t^8 + 44299317t^{10} + \mathcal{O}(t^{12})$	$32t^2 + 2592t^4 + 97984t^6 + 2366496t^8 + 41557344t^{10} + \mathcal{O}(t^{12})$	$1 + 123t^2 + 6710t^4 + 220812t^6 + 5027328t^8 + 85856661t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{5',6} _{N=2}$	$1 + 183t^2 + 16626t^4 + 1000427t^6 + \mathcal{O}(t^8)$	$128t^3 + 21632t^5 + \mathcal{O}(t^7)$	$1 + 183t^2 + 128t^3 + 16626t^4 + 21632t^5 + 1000427t^6 + \mathcal{O}(t^8)$

$\text{MQ}_{6,6} _{N=1}$	$1 + 92t^2 + 5696t^4 + 235812t^6 + 6925608t^8 + 153422070t^{10} + \mathcal{O}(t^{12})$	$64t^2 + 5248t^4 + 230784t^6 + 6881472t^8 + 153102720t^{10} + \mathcal{O}(t^{12})$	$1 + 156t^2 + 10944t^4 + 466596t^6 + 13807080t^8 + 306524790t^{10} + \mathcal{O}(t^{12})$
$\text{MQ}_{7,7} _{N=1}$	$1 + 138t^2 + 16303t^4 + \mathcal{O}(t^6)$	$128t^2 + 16128t^4 + \mathcal{O}(t^6)$	$1 + 266t^2 + 32431t^4 + \mathcal{O}(t^6)$

D Central charges and 3d mirrors for D-type class-S theories

In this appendix we briefly review two properties of class-S theories of D_N type. We restrict ourself to the case in which the 6d $(2,0)$ D_N theory is compactified on the sphere, with regular (untwisted) punctures only. We discuss three instances of theories of this type in the main text, in sections 3.10, 3.11 and 3.12. There, we consider examples in which a single three-punctured sphere describes the 4-dimensional version of the $E_n \times E_n$ theory ($n = 6, 7, 8$), i.e. two copies of E_n Minahan-Nemeschansky.

The first property that we review in this appendix is the rule for the computation of superconformal central charges a and c , giving the data labeling the punctures [26]. The second property is the prescription for finding the corresponding 3d $\mathcal{N} = 4$ mirror theories [18].

D.1 Central charges

The central charges a and c of a 4d $\mathcal{N} = 2$ SCFT are defined via the trace anomaly in a curved background,

$$T_\mu{}^\mu = \frac{c}{16\pi^2}(\text{Weyl})^2 - \frac{a}{16\pi^2}(\text{Euler})^2. \quad (\text{D.1})$$

For Lagrangian theories, the central charges are related to the number of vector multiples and hypermultiplets by

$$a = \frac{5n_v + n_h}{24}, \quad c = \frac{2n_v + n_h}{12}. \quad (\text{D.2})$$

For non-Lagrangian theories, formula (D.2) still holds, but now n_h and n_v are interpreted as an effective number of hypermultiplets and vectormultiplets. For the subset of theories of our current interest, it holds that

$$\begin{aligned} n_v &= -\frac{1}{3}N(16N^2 - 24N + 11) + \sum_\alpha \delta n_v^{(\alpha)}, \\ n_h &= -\frac{8}{3}N(N-1)(2N-1) + \sum_\alpha \delta n_h^{(\alpha)}, \end{aligned} \quad (\text{D.3})$$

where g is the genus of the Riemann surface and α runs over the set of punctures. $\delta n_v^{(\alpha)}$ and $\delta n_h^{(\alpha)}$ are local contributions coming from the α -th puncture. Both formulae (D.3) and an algorithmic rule for the computation of δn_h and δn_v are derived in [26]. In the same paper, the explicit values of δn_h and δn_v is listed for all the punctures up to $N = 6$. We defer the reader to such paper for further details.

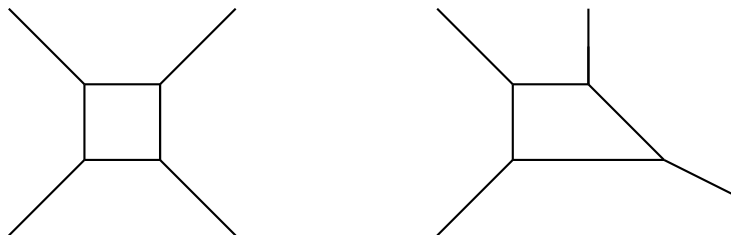


Figure 8. 5-brane web diagrams for pure $SU(2)_0$ (left) and pure $SU(2)_\pi$ (right) gauge theories.

D.2 3d mirrors

The procedure consists in associating an orthosymplectic linear quiver to each puncture of the four dimensional theory. For 4d theories of D_N type, each of those 3d $\mathcal{N} = 4$ quiver tails ends with a flavor symmetry node $SO(2N)$. As a second step, one glues the flavour nodes together, producing a star-shaped quiver.

The rule for associating the quiver tails to the punctures is the following. Consider a regular puncture whose Nahm partition is given by $[h_1, \dots, h_J]$. The associated 3d quiver tail is

$$[SO(2N)] - \text{USp}(r_1) - SO(r_2) - \dots - \text{USp}(r_{J-1}), \quad (\text{D.4})$$

where the quantities r_a are defined as

$$r_a = \left[\sum_{b=a+1}^J h_b \right]_{+,-}, \quad + : SO, \quad - : \text{USp} \quad (\text{D.5})$$

and $[n]_{+(-)}$ is the smallest (resp. largest) even integer $\geq n$ (resp. $\leq n$). When $r_{J-1} = 0$, we remove the last group $\text{USp}(0)$. We defer the reader to [18] for further details.

E 5-brane webs for $SO(4)$ theory with different discrete theta angles

It is known that the pure $SU(2)$ gauge theory has two choices, $SU(2)_0$ and $SU(2)_\pi$, depending on their discrete theta angles. $SU(2)_0$ and $SU(2)_\pi$ gauge theories are known to have different UV fixed points with different global symmetry, $E_1 = SU(2)$ and $\tilde{E}_1 = U(1)$, respectively. They also have distinct 5-brane configurations as depicted in figure 8. One can readily see from the 5-brane webs given in figure 8 that $SU(2)_0$ gauge theory has an one-dimensional Higgs branch at infinite coupling, while $SU(2)_\pi$ gauge theory has no Higgs branch at infinite coupling.

In this appendix, we discuss the 5-brane configurations for $SO(4) = SU(2) \times SU(2)$ gauge theory with different discrete theta angles. As each $SU(2)$ can have discrete theta angle, we introduce the following shorthand notation to denote two discrete theta angles for $SO(4)$, $SO(4)_{\theta_1, \theta_2} = SU(2)_{\theta_1} \times SU(2)_{\theta_2}$. It follows that a conventional 5-brane web for $SO(4)$ given in figure 9 corresponds to a 5-brane web for $SO(4)_{0,0}$. At the infinite coupling, this 5-brane web for $SO(4)_{0,0}$ is deformed to figure 1, from which one can read off the magnetic quiver and the corresponding Hilbert series. In particular, we can check that the Higgs branch dimension is two, as expected. We also find the $SU(2)$ global symmetry from the parallel

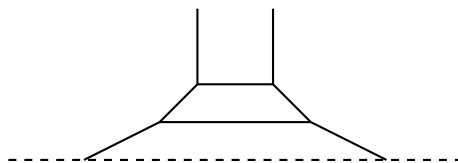


Figure 9. 5-brane web diagram for $SO(4)$.

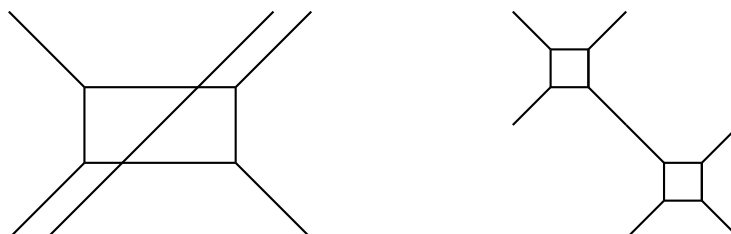


Figure 10. 5-brane web for $SU(2)_0$ - $SU(2)_0$ quiver. To read off the theta angles we deform the web on the left to that on the right.

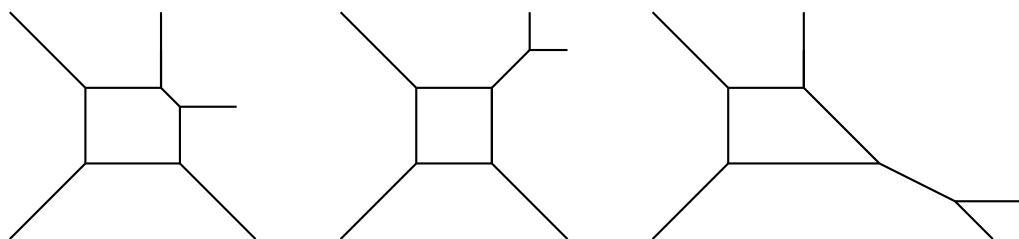


Figure 11. Left: 5-brane web diagram for $SU(2) + 1\mathbf{F}$. Middle: decoupling the hypermultiplet by giving it a positive mass. Taking the mass to infinity results in the web diagram for pure $SU(2)_0$ theory. Right: decoupling the hypermultiplet by giving it a negative mass. Taking the mass to negative infinity results in the web diagram for pure $SU(2)_\pi$ theory.

NS5-branes, which is part of the $E_1 \times E_1$ symmetry. The $SU(2)_0 \times SU(2)_0$ gauge theory can also be constructed from the 5-brane web without an O5-plane as given in figure 10. We find that if the bifundamental mass is large enough, this web can be decomposed into two copies of $SU(2)_0$ webs.

To construct 5-brane webs for $SO(4)_{\theta_1, \theta_2}$ gauge theories with different discrete theta angles, we first recall how 5-brane web for $SU(2)_0$ and $SU(2)_\pi$ can be obtained from that for $SU(2)$ gauge theory with one flavor. Pure $SU(2)_0$ theory is obtained by decoupling the flavor with a positive infinity mass, while pure $SU(2)_\pi$ theory is obtained by giving a negative infinity mass, as depicted in figure 11.

We then consider an analogous situation with 5-brane web with an O5-plane. For instance, consider a 5-brane web with spinor matter as in figure 12. The left of figure 12 is the web for $SO(4)$ gauge theory with one spinor and one conjugate spinor, which corresponds to two copies of $SU(2)$ gauge theories with one flavor. If we take all of their masses to be positive infinity, which is to move the branes corresponding to the spinor matters to the right, then the resulting 5-brane web trivially goes back to the pure $SO(4)_{0,0}$ gauge theory in figure 9. On the other hand, one may instead bring the spinor matter closer

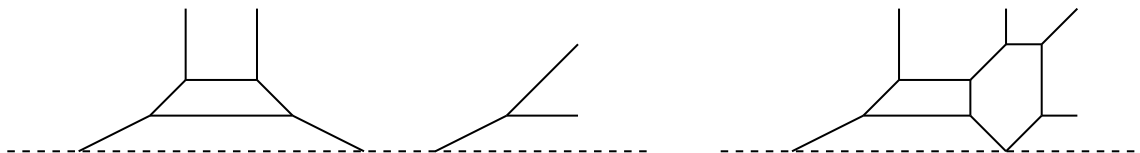


Figure 12. Left: the web diagram for $\text{SO}(4)$ gauge theory with $1\mathbf{s} + 1\mathbf{c}$. Right: the web after flop transitions.

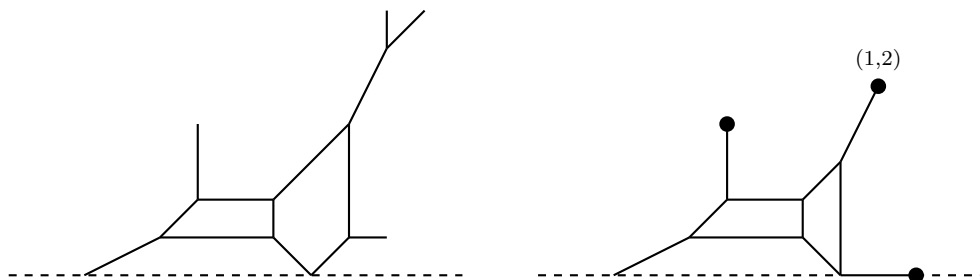


Figure 13. Left: the web diagram for $\text{SO}(4)$ gauge theory with $1\mathbf{s} + 1\mathbf{c}$ with negative large masses. Right: 5-brane web diagram for $\text{SO}(4)_{\pi,\pi}$.

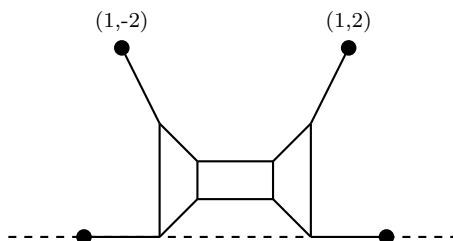


Figure 14. 5-brane web diagram for $\text{SO}(4)_{2\pi,2\pi}$.

to the pure $\text{SO}(4)_{0,0}$ followed by a sequence of flop transitions including the “generalized flop transition” [40], which yields the right of figure 12. When we consider their masses to be negatively large enough, further flop transition is induced to yield the left of figure 13. Taking their masses to be identical and further tuning this mass to be negative infinity, we find the $\text{SO}(4)_{\pi,\pi}$ gauge theory in the right of figure 13. We can check from this web that there is no Higgs branch. Also, we find no non-Abelian global symmetry, as expected.

We can even construct a 5-brane web for $\text{SO}(4)_{2\pi,2\pi}$ gauge theory, by repeating the procedure above. Given a 5-brane web for $\text{SO}(4)_{\pi,\pi}$ as in figure 13, we can introduce another set of spinor and conjugate spinor on the left. By taking their masses to be negatively large, we obtain the 5-brane web in figure 14. This is a 5-brane web for $\text{SO}(4)_{2\pi,2\pi}$ gauge theory, whose Higgs branch should be identical to $\text{SO}(4)_{0,0}$ gauge theory. From this 5-brane web diagram, we can explicitly check using [2] that it has Higgs branch with dimension two at the infinite coupling, as expected. We note that it is also interesting to observe that the S-dual of this web diagram looks like $\text{SU}(2) \times \text{SU}(2)$ gauge theory as a special case of the D-type Dynkin quiver gauge theory constructed with ON^0 -plane, where the gauge coupling constants are tuned to be identical. The symmetric shape indicates that the corresponding $\text{SU}(2)$ theories have discrete theta angle 0 rather than π .

All these web diagrams for $SO(4)_{\theta_1, \theta_2}$ gauge theories with different discrete theta angles and their consistency checks give further support for the correspondence between (3.1) and (3.2).

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